Investigation of angstrom-thick aluminium oxide passivation layers to improve the gate lag performance of GaN HEMTs

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Keywords: AlGaN, atomic layer deposition, dielectric, GaN, gate lag, HEMT, passivation

Abstract
In this paper, we report an angstrom-thick atomic layer deposited (ALD) aluminum oxide (Al₂O₃) dielectric passivation layer for an AlGaN/GaN high electron mobility transistor (HEMT). Our results show a 55% improvement in the gate lag performance of the design and a decrease by half in interface state density upon coating with two cycles of ALD Al₂O₃. DC characteristics such as current density, threshold voltage, and leakage currents were maintained. ALD Al₂O₃ passivation layers with thicknesses up to 10 nm were investigated. XPS analyses reveal that the first ALD cycles are sufficient to passivate GaN surface traps. This study demonstrates that efficient passivation can be achieved in atomic-scale with dimensions much thinner than commonly used bulk layers.

1. Introduction

GaN based high electron mobility transistors (HEMTs) are advantageous for high power and high frequency applications due to their superior properties such as wide band gap, high breakdown field, and high saturation velocity [1–3]. Although GaN HEMTs are widely preferred for a variety of commercial and military applications, unpassivated devices suffer from limitations in characteristics such as radio frequency (RF) current, transconductance, and breakdown voltage as a result of the trapping of electrons in the surface states created by defects and dislocations [4, 5]. Surface passivation has become the most established approach to address the problem of current collapse, due to its simplicity and efficacy, with SiNx being the most commonly preferred dielectric [6–8]. SiNx passivation schemes are particularly effective for III-N devices due to the relatively low state density at the SiNx/III-N interface and the possibility of in situ MOCVD deposition [9]. Numerous other dielectrics have been demonstrated as passivation layers over the years with different advantages [10–12]. Al₂O₃ passivation has been demonstrated with record drain current and transconductance and improved pulsed I-V characteristics [13–15]. Al₂O₃ is an attractive material because of its large bandgap, relatively high dielectric constant, and high breakdown field and has the advantage of high quality oxide/III-N interfaces when atomic layer deposited (ALD) [16].

Recently, angstrom-thick passivation has been demonstrated as an efficient route to provide both surface coverage and trap passivation in various applications [17–20]. It was found that the first ALD cycles will only passivate surface defects without introducing any new bulk traps. Thus, it is envisioned that this atomic-scale selective passivation is the optimal way for GaN HEMTs.

In this paper, we demonstrate that two ALD cycles of Al₂O₃ will lead to about 55% improvement in the gate lag characteristic of the GaN HEMT compared to that of an unpassivated design. This improvement has been correlated with the reduction in the density of GaN surface traps, which has been studied with frequency dependent conductance measurements. An in-depth XPS surface analysis confirmed that 2 Å Al₂O₃ is sufficient
for passivation of the AlGaN surface. DC characteristics such as knee voltage and transconductance were improved and other characteristics were maintained for the passivated device.

2. Device fabrication

The epitaxial structure consists of a 20 nm AlN nucleation layer and 1350 nm Carbon doped high resistive buffer layer on a semi-insulating SiC wafer, followed by 150 nm high mobility GaN channel, 1 nm AlN spacer, and 22 nm AlGaN barrier layer with 26% Aluminum content capped with 3 nm GaN. The electron mobility and 2DEG density were found to be 2061 cm$^2$ V$^{-1}$ s and 1.07 $\times$ 10$^{13}$ cm$^{-2}$ respectively. Average sheet resistivity was 284 Ohms/Sq across the 3 inch wafer. Device fabrication began with mesa isolation. Ohmic contacts were formed with electron beam deposited Ti/Al/Ni/Au stack, followed by 30 s rapid thermal annealing at 850 °C in N$_2$ ambient. 250 nm gate contacts were formed with electron beam lithography and Ni/Au gate metal deposition. For device passivation Al$_2$O$_3$ depositions were carried out at 200 °C in an ALD reactor (Cambridge Nanotech Savannah S100) employing Al(CH$_3$)$_3$ solution as the deposition precursor. The pulse and purge durations were 0.015 and 10 s, respectively. Water was used as the oxygen precursor. The deposition rate was found as 1 Å/cycle. Different devices were fabricated with 2 cycles, 4 cycles, 10 cycles, and 100 cycles to obtain passivation layer thicknesses of 2 Å, 4 Å, 1 nm, and 10 nm, respectively. The devices in this letter have a gate width of 2 $\times$ 125 μm, a gate-source spacing of 1.5 μm, and gate drain spacing of 3.25 μm. Schottky diodes were fabricated with a similar process flow for capacitance measurements.

3. Results and discussion

Capacitance-voltage (C-V) measurements of fabricated Schottky diodes were conducted in order to examine the gate characteristics of the different passivation schemes. Figure 1 shows the C-V characteristics for the unpassivated, 2 Å, 4 Å, and 10 nm ALD Al$_2$O$_3$ passivated samples measured at a frequency of 1 MHz. For the unpassivated, 2 Å, and 4 Å passivated samples similar capacitance curves are obtained. An increase in capacitance is observed for the 10 nm Al$_2$O$_3$ passivated sample, confirming the formation of the oxide layer. A negative threshold shift is also observed for the sample with 10 nm passivation.

Frequency dependent conductance measurements were carried out to determine the interface state density $D_{it}$. The frequency was varied from 10 kHz to 2 MHz. The parallel $C_m$—$G_m$ circuit was considered for the equivalent circuit to the interface. The interface state density was obtained from [21] as

$$G_p = \frac{q\omega\tau_m D_{it}}{1 + (\omega\tau_m)^2} = \frac{\omega G_m C_{ox}^2}{C_m + \omega^2(C_{ox} - C_m)^2}$$

where $C_m$ and $G_m$ are the measured capacitance and conductance at frequency $\omega$, and $C_{ox}$ is the oxide capacitance. The calculated $G_p/\omega$ values for the unpassivated, 2 Å, 4 Å, and 10 nm ALD Al$_2$O$_3$ passivated schemes are shown in Figure 2. The extracted $D_{it}$ values are given in Table 1. The relative position of the trap energy level
with respect to the conduction band edge is found to be in the range of 0.18–0.19 eV and the time constants in the range of 0.19–0.32 μs for the samples. An interface state density of $3.08 \times 10^{13}$ eV$^{-1}$ cm$^{-2}$ is obtained for the unpassivated sample. For the 2 Å Al$_2$O$_3$ sample the lowest $D_{it}$ of the studied samples is obtained, with a value of $1.62 \times 10^{13}$ eV$^{-1}$ cm$^{-2}$, which corresponds to nearly a 50% decrease in $D_{it}$ compared to the unpassivated sample. For the 4 Å Al$_2$O$_3$ sample, however, the highest $D_{it}$ of $2.36 \times 10^{13}$ eV$^{-1}$ cm$^{-2}$ is calculated, which corresponds to a nearly 50% increase compared to the unpassivated sample. Thus, the conductance method measurements indicate that angstrom thick ALD passivation layers are able to passivate surface states selectively whereas thicker layers lead to increased surface state density.

Table 1 lists the DC measurement results and calculated gate lag percentages for different devices, that are, unpassivated, 2 Å, 4 Å, 1 nm, and 10 nm ALD deposited Al$_2$O$_3$. Figure 3 compares the transfer characteristics and gate currents of the devices for each passivation scheme, at $V_D = 10$ V. The threshold voltage was obtained as −4.06 V for the unpassivated device using linear extrapolation at maximum linear slope and does not vary significantly for 2 Å, 4 Å, and 1 nm Al$_2$O$_3$ passivation. For 10 nm ALD Al$_2$O$_3$ passivation, $V_{th}$ decreases to −4.15 V, analogous to the decrease observed in the C-V measurement, due to the increased sheet charge from the passivation layer. The measured output characteristics are shown in Figure 4. The drain current density $I_{ds}$ (at a gate voltage of $V_G = 1$ V) was obtained as 0.75 A mm$^{-1}$ for the unpassivated reference device, which increased slightly by 5% to a value of 0.79 A mm$^{-1}$ for the 4 Å and 1 nm passivation scheme and by 10% to a maximum value of 0.82 A mm$^{-1}$ for the 10 nm passivation scheme. Similarly, the peak transconductance value increased by 4.7% from 179 mS mm$^{-1}$ for the reference device to 188 mS mm$^{-1}$ for the devices with 1 nm and 10 nm passivation. An improvement in knee voltage is observed with increasing Al$_2$O$_3$ passivation thickness. The reason for these improvements in the drain current characteristics is related to enhanced electron transport caused by an increase in sheet carrier concentration, due to the reduction in surface states. The drain leakage current $I_{leak}$ is extracted from the transfer characteristics measurements as the current at $V_G = −6$ V and $V_D = 10$ V. $I_{leak}$ increases monotonically with passivation layer thickness, which is also

![Figure 2](image-url). Calculated values of trap conductance as a function of angular frequency for unpassivated, 2 Å, 4 Å, and 10 nm ALD Al$_2$O$_3$ passivated AlGaN/GaN HEMTs.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Units</th>
<th>Unpassivated</th>
<th>2 Å</th>
<th>4 Å</th>
<th>1 nm</th>
<th>10 nm</th>
</tr>
</thead>
<tbody>
<tr>
<td>$D_{it}$</td>
<td>eV$^{-1}$ cm$^{-2}$</td>
<td>$3.08 \times 10^{13}$</td>
<td>$1.62 \times 10^{13}$</td>
<td>$2.36 \times 10^{13}$</td>
<td>—</td>
<td>$6.52 \times 10^{13}$</td>
</tr>
<tr>
<td>$I_{ds}$</td>
<td>A mm$^{-1}$</td>
<td>0.75</td>
<td>0.78</td>
<td>0.79</td>
<td>0.79</td>
<td>0.82</td>
</tr>
<tr>
<td>$V_{th}$</td>
<td>V</td>
<td>−4.06</td>
<td>−4.07</td>
<td>−4.07</td>
<td>−4.07</td>
<td>−4.15</td>
</tr>
<tr>
<td>$g_{m}$</td>
<td>mS mm$^{-1}$</td>
<td>179</td>
<td>186</td>
<td>187</td>
<td>188</td>
<td>188</td>
</tr>
<tr>
<td>$I_{leak}$</td>
<td>nA mm$^{-1}$</td>
<td>2.7</td>
<td>3.4</td>
<td>1.9</td>
<td>1.6</td>
<td>21.4</td>
</tr>
<tr>
<td>$I_{leak}$</td>
<td>nA mm$^{-1}$</td>
<td>8.3</td>
<td>12.7</td>
<td>19.3</td>
<td>20.3</td>
<td>206</td>
</tr>
<tr>
<td>Gate Lag</td>
<td>%</td>
<td>20.3</td>
<td>9.4</td>
<td>11.0</td>
<td>11.8</td>
<td>8.5</td>
</tr>
</tbody>
</table>
attributed to the increasing sheet charge of the 2DEG. The Schottky characteristics of the AlGaN/GaN HEMTs were measured with a floating drain contact (Figure 5) and the gate leakage current $I_{g,\text{leak}}$ is obtained from this measurement as the gate current at $V_{G} = -6$ V. The reverse gate leakage current, in contrast to the drain leakage current, does not follow a monotonically increasing trend, indicating that the reverse gate leakage current is not dominated by surface leakage. The reverse gate leakage current increases for the initial passivation 2 Å, decreases for the two subsequent passivation thicknesses, then increases to its highest value for the greatest passivation thickness of 10 nm. The dominant mechanism of gate leakage is attributed to be edge gate leakage current, as the trend of the reverse gate leakage current for increasing passivation thickness correlates inversely with the trend in the gate lag characteristic. A decrease in gate lag corresponds to an increase in gate leakage, indicating that extension of the gate depletion region towards the drain by the virtual gate effect acting to disperse the electric field lines induced by the depletion region reduces the edge current [22]. Passivation that suppresses the virtual gate thereby leads to an increase in gate leakage current. Low forward gate leakage current and nearly constant gate turn on voltage is maintained for each passivation scheme.

To further understand the correlation between the passivation layer and the surface states gate lag measurements were carried out. The gate lag measurements were performed using an Agilent E3631A power supply, Keysight Technologies 33500B waveform generator, and a Keysight InfiniVision DSOX2004A
oscilloscope. The measurements were carried out by pulsing the gate from \(-6\) V to 0 V, with a pulse width of 1 \(\mu\)s and period of 20 \(\mu\)s with the drain kept at 0 V. The pulsed drain current response with respect to the DC values are compared in Figure 6 for the studied passivation schemes. A gate lag of 20.9% is obtained for the unpassivated device, which improves by 55% to a gate lag value of 9.5% for the device with 2 Å passivation. For the 4 Å and 1 nm \(\text{Al}_2\text{O}_3\) passivation layers, a slight increase in the gate lag characteristics is observed, however, significant improvement compared to the unpassivated case is maintained. For 10 nm \(\text{Al}_2\text{O}_3\) passivation layer the lowest gate lag of 8.7% is observed. 

Figure 5. \(I_C-V_G\) characteristics measured with floating drain contact for unpassivated, 2 Å, 4 Å, 1 nm, and 10 nm ALD \(\text{Al}_2\text{O}_3\) passivated AlGaN/GaN HEMTs.

Figure 6. Drain current response of the pulsed gate voltage normalized to DC values with \(V_D = 0\) V and \(V_G\) pulsed from \(-6\) V to 0 V for unpassivated, 2 Å, 4 Å, 1 nm, and 10 nm ALD \(\text{Al}_2\text{O}_3\) passivated AlGaN/GaN HEMTs. The pulsewidth is 1 \(\mu\)s at 5% duty cycle.

To gain insight on the passivation property of ALD \(\text{Al}_2\text{O}_3\) layer, x-ray photoelectron spectroscopy (XPS) is utilized for unpassivated and 2 Å \(\text{Al}_2\text{O}_3\) coated samples. Before ALD process, all the samples have been treated with a diluted HCl acid to remove the natural oxide layer. Figure 7 shows the N1s and Ga3d spectra of these two samples. Peak position correction was calibrated by referencing the C1s orbital peak position (284.8 eV) and the other peaks in the spectrum were shifted accordingly. As previously explained in several reports, the dominant surface defects of GaN are Ga and N vacancies (or dangling bonds). Based on the calculation of free energy by classical nucleation theory, most of the oxygen-derived hydroxyl groups such as OH radicals and \(\text{H}_2\text{O}\) or \(\text{O}_2\) will
be chemisorbed near imperfections such as dangling bonds and vacancies. Looking back into Figure 7(a), the N1s spectra is deconvoluted into five main peaks \[24\]; three of which are assigned into Auger Ga LMM peaks, the dominant one comes from the Ga-N bond, and the one in the higher energy tail is attributed to N-O bonds. As we can see, the passivated sample has a stronger N-O related peak. During the ALD process, the Ga vacancy positions are passivated by oxygen molecules and form these N-O bonds. On the other side, the Ga3d spectra can be scrutinized to gain a further insight on the surface properties of the GaN layer. As illustrated in Figure 7(b), this spectrum is deconvoluted into three Gaussian profiles \[25\]; a broad and weak response originated from N2s orbitals, a dominant peak assigned to Ga-N bond, and a high energy response from Ga-O bonds. Looking at these peaks, it can be clearly seen that the Ga-O related peak is more dominant for 2 Å Al2O3 passivated case (compared to the unpassivated one). Therefore, similar to N1s data, the Al2O3 passivation layer will substitute N vacant positions and facilitate formation of Ga-O bonds. All the above-mentioned results confirm the efficient passivation of surface traps upon coating with the 2 Å Al2O3 layer. Oxygen-containing gas molecules tend to be chemisorbed on the surface of a semiconductor host through the capture of free electrons. Consequently, these chemisorbed radicals reduce the density of free carriers in the vicinity of the semiconductor surface and deplete the surface electron states. This, in turn, triggers the existence of the space charge region and induces band bending near the interface. On the other hand, these chemisorbed oxygen molecules are likely attached into trap states and dangling bonds. The passivation of surface traps, in first ALD cycles, reduces the surface traps density, mitigates adsorption of oxygen radicals, and consequently reduces band bending. However, as we go to larger ALD cycles, bulk trap states in Al2O3 layer start to become dominant and therefore diminish the abovementioned characteristics. Moreover, the Al2p spectra have been extracted for 2 Å and 10 nm coated samples. As Figure 8 clearly illustrates, the portion related to Al element is the major peak for the thin passivated sample. However, as we increase the Al2O3 layer thickness to 10 nm, the spectrum is mainly attributed to the
oxide related peak. Thus, the first two ALD cycle just passivates the defect states and does not form a continuous Al$_2$O$_3$ layer. However, the subsequent cycles trigger the formation of a continuous layer.

4. Conclusion

In conclusion, an angstrom-thick ALD Al$_2$O$_3$ dielectric is reported for the passivation of surface traps in AlGaN/GaN HEMTs and passivation thicknesses up to 10 nm are studied. For the a passivation layer as thin as 2 Å Al$_2$O$_3$, a greater than 50% improvement in gate lag compared to the unpassivated device is achieved. Improvements in DC characteristics such as drain current density, knee voltage, and transconductance are observed. XPS analysis confirms that Ga vacancy positions are are passivated by oxygen molecules for Al$_2$O$_3$ as thin as 2 Å.

Acknowledgments

This work is supported by the TUBITAK under Project No. 116F041. One of the authors (E.O.) also acknowledges partial support from the Turkish Academy of Sciences.

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Figure 8. Al2p high resolution XPS pattern of the resulting 2 Å thick and 10 nm Al2O3 coated samples.


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