

## RESONANCE LENS ANTENNA ANALYSIS FOR MM-WAVE APPLICATIONS

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We report what is to our knowledge the first accurate theoretical investigation of the electromagnetic behavior of 2-D elliptical lenses of finite wavelength-scale size. The role of internal resonances in the focal domain formation is studied. A proposal of a narrow-band receiver based on a hemielliptic lens tuned to a resonance is discussed. Possible features of such a lens-coupled receiver are stability of the resonance field with respect to the angle of arrival of incident wave and several times greater values of the peak field intensity that may potentially lead to higher sensitivity and better scanning performance.

In the analysis, we use the Muller boundary integral equation (BIE) technique. This full-wave mathematically rigorous method is combined with trigonometric Galerkin discretization to result in the efficient numerical solution for an arbitrary set of the electrical, geometrical, and material parameters.

Numerical results are generated for a quartz elliptical lens ( $\epsilon = 3.8$ ) with dimensions typical to mm-wave radar applications [1]. Near field analysis, lens-focusing properties and lens frequency-dependent performance are presented.

### Introduction.

Printed slot or strip elements combined with dielectric lenses are frequently met components of mm and sub-mm wave receivers [2]-[6]. The attention they have been attracting recently is due to their capability of compact integration with other electronic components such as detecting diodes, local oscillators and mixers. Besides, they provide better efficiency than other types of antennas printed on homogeneous substrates. The elliptical shape of the lens is borrowed from optics to provide focusing properties if its eccentricity is selected properly. On the other hand, the lens interface gives rise to reflections inside a realistic lens that may significantly affect the input impedance and sensitivity. This aspect, which has not been properly investigated in the literature, is a critical point in the overall design of lens antennas [3], [6]. Various analytical techniques used to simulate dielectric lenses have been commonly based on the ray-tracing, neglected the lens size and curvature and hence failed to characterize the internal resonances. The validity of such approximations can in many cases be questionable as the actual size of the lens is several wavelengths [2]-[7]. In fact, a behavior of such a lens should clearly display both ray-like and resonance-mode features. On the other hand, direct numerical simulators like FDTD-based ones are extremely time and memory consuming, require

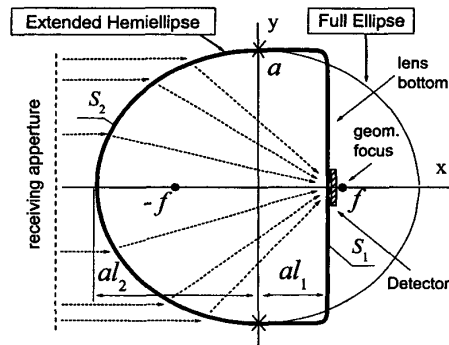


Figure 1. Geometry and notations of an extended hemielliptic lens. Dotted arrows are for the ray-tracing focusing diagram.

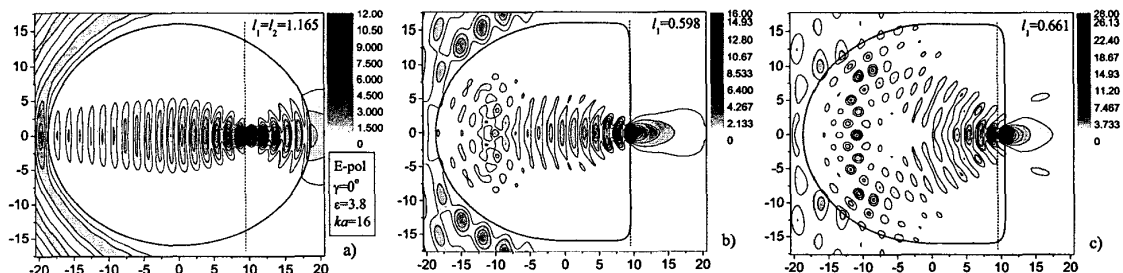


Figure 2. The near-field intensity portraits for the elliptical lens (a) and extended hemielliptic lenses with  $l_1$  corresponding to the geometrical focus location (b) and the nearest resonance (c). Dotted lines denote the far geometrical focus location.  $E$ -polarization,  $\gamma = 0$ ,  $\epsilon = 3.8$ ,  $ka = 16$ ,  $l_2 = 1,165$ .

accurate treatment of the back-reflections from the edge of the computational window, and may suffer from a loss of accuracy. All the above shows that the true electro-magnetic behavior of dielectric lenses is still far from being clear.

**Outline of the solution.**

A brief outline of analysis method is as follows. In 2-D, the focusing properties study can be formulated as a problem of the plane wave scattering by a smooth dielectric cylinder. The application of the Green's identity enables one to express the fields inside and outside the lens versus the values of fields and their normal derivatives at cylinder's contour. By applying the boundary conditions, a set of the Fredholm second kind IEs is obtained, known as Muller's boundary IEs [8]. Further we discretize them by applying the Galerkin method with entire-domain angular exponents as expansion functions. To treat the (integrable) singularities in some of the kernel functions, we use analytical integration of the canonical circular-cylinder counterpart terms.

This procedure results in a Fredholm second kind infinite-matrix equation having favorable features, with matrix elements and right-hand-part terms obtained as Fourier-expansion coefficients of some twice-continuous functions. They can be economically computed by using the DFFT and FFT algorithms, respectively. Developed algorithm guarantees point-wise convergence of the numerical solution, i.e., a possibility to minimize the error to the level set by the DFFT if solving progressively greater matrices. Note that required accuracy can be achieved for an arbitrary set of the lens parameters, e.g., for any value of the contrast between the lens material and the background medium. It is also free from the inaccuracies near to the sharp natural resonances that is intrinsic to conventional numerical approximations [9], as well as from the well-known defect of "numerical resonances" occurring if the fields are presented as single or double layer potentials [10].

More solution details and algorithm properties can be found in [11]. The advantages of the method outlined above make it efficient CAD tool, which can be applied to the accurate analysis of dielectric lenses in the wide range of frequencies and geometries.

**Numerical results.**

We shall simulate the cross-sectional contour of the truncated elliptic lens with a curve  $S$ , which is twice-differentiable at every point and consists of two open curves  $S_1, S_2$  smoothly joined together at the points marked with crosses (Fig.1). Here,  $S_2$  is a half of the ellipse with eccentricity  $e=1/\sqrt{\epsilon}$  and  $S_1$  is a half of so-called "super-ellipse" that simulates a rectangle with rounded corners.

In the Fig. 2 the near-field intensity portraits for the extended hemielliptical lens with  $l_1$  corresponding to the geometrical focus location and the first resonance revealed in Fig. 3 are shown. It is seen that the ellipse truncation leads to a dramatic change in the field distribution inside the lens and a noticeable increase of the peak intensity. Moreover, while the field intensity portrait for the lens extended to the geometrical focus resemble a ray-tracing picture (here and after quasi-optical lens), similar portrait for the lens extended to the nearest internal resonance

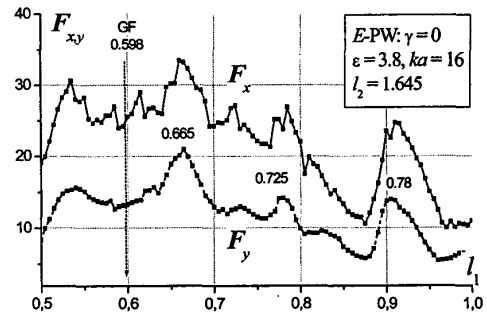


Figure 3. Focusability of the lens along x- and y-axis versus the parameter characterizing the extension of the lens rear side.

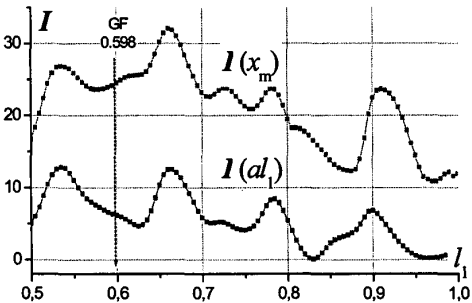


Figure 4. The overall peak field intensity and that on the lens bottom versus the rear side extension parameter.

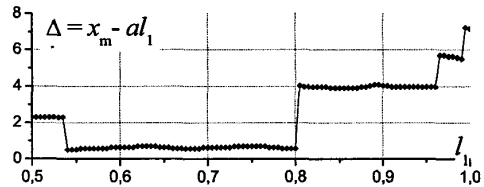


Figure 5. The distance between the point with the highest intensity and the lens bottom.

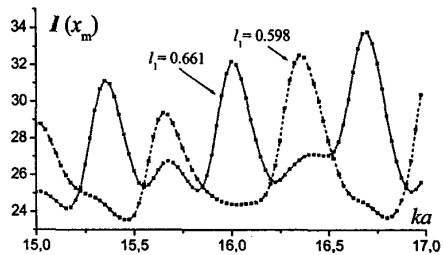


Figure 6. The peak field intensity versus the normalized frequency parameter for the quasi-optical and the resonant lenses.

(here and after resonant lens) clearly shows additional intensive field spots. Such resonance resembles, to some extent, asymptotic GO resonances known as "billiard-type" ones. Here they have specific triangular portraits and can be classified by two indices, which indicate the variations of the field along the lens flat side ( $m$ ) and a side of the triangle ( $n$ ). Two neighboring resonances in  $l_1$  have  $n$  and  $n+1$  variations with apparently the same  $m$ , and the distance between them corresponds to half a wavelength in dielectric along the triangle side.

The ability of the lens to focus the radiation along the  $x$ - and  $y$ -axes can be quantified by the "focusability parameters" defined as  $F_z = aI(x_m, y_m)/|z_m - z_{0.5}|$ , where  $z = x, y$ ,  $(x_m, y_m)$  are the coordinates of the highest field intensity point ( $I = |E_z|^2$ ) inside the lens, and  $(x_{0.5}, y_{0.5})$  are the coordinates of the half maximum intensity values.

Analysis of graphs presented in Fig. 3 and Fig. 4 shows that the focal spot size is stable for the considered values of  $l_1$ . Internal resonances result in both focusability and peak intensity increase. Fig. 5 shows that the focal domain is always shifted inside the lens but located at the fixed distance from the lens bottom for a wide range of values of  $l_1$ . Besides, as it is seen from Fig. 4, the field intensity at the lens bottom is only a half its value in focal point. These facts are important as they suggest one to locate the detectors inside the lens instead of typical lens-coupled antenna design in which a lens is glued over a slot antenna.

Analysis shows that such a resonant lens tuned for the internal resonance has certain advanced properties, despite losing its wideband nature, namely stability of the resonance field with respect to the angle of arrival of incident wave and several times greater values of the peak intensity. Narrow-band receiver exploiting this effect may potentially have improved sensitivity and scanning performance.

The frequency dependences of the peak intensity for quasi-optical and resonant lenses (Fig. 6) show that both types of lenses have internal resonances, which result in narrow peak intensity flashes. This fact has two outcomes. One is that any finite size lens focusing properties are frequency-dependent. Another is that for a wide-band applications the only requirement for the lens extension size is to be in a range of 0.55-0.8 in order to keep the focal domain close to the lens bottom because the presence of internal resonances will similarly affect any lens focusing characteristics with a bottom extension in the given range.

#### Summary.

The electromagnetic behavior of extended hemi-elliptical lens has been studied. The most important feature revealed by the accurate analysis is that the resonances may play a dominant role in the wavelength-scale lens behavior. Obtained results show that advanced characteristics of the lens-coupled receiver can be achieved, although in a narrow band, by exploiting a resonance.

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