



# An All-Dielectric Metasurface Coupled with Two-Dimensional Semiconductors for Thermally Tunable Ultra-narrowband Light Absorption

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## Abstract

Two-dimensional (2D) transition metal dichalcogenides (TMDCs) have attracted tremendous attention over the past decades. Due to their unique features such as high mobility and direct bandgap, they are suitable candidate for the optoelectronic devices. However, due to their ultrathin thickness, their optical absorption is quite weak, and therefore, a trapping scheme for strong light-matter interaction is essential to overcome this deficiency. To accomplish strong light absorption, loss-less dielectric-based metasurfaces with ideally no parasitic absorption are excellent choices. Herein, we report an ultra-narrowband thermally tunable all-dielectric metasurface coupled absorber with TMD monolayer. In this proposed structure, high absorption with ultra-narrow full-width-at-half-maximum (FWHM) is achieved. Different design configurations are studied to find the most suitable structure. In the optimized design, an absorptance as high as 0.85 with a FWHM of 3.1 nm is achieved. This structure also shows thermal sensitivity of 0.0096 nm/°C, without the use of any phase change material component. This architecture can be used as a 2D and highly efficient tunable single-color photodetector. The proposed dielectric metasurface can be adopted for other types of 2D and ultrathin semiconductor-based optoelectronics.

**Keywords** Photodetector · Metasurface · Two-dimensional material · Titanium dioxide

## Introduction

In recent years, two-dimensional (2D) materials such as graphene, hexagonal boron nitride, and transition metal dichalcogenides (TMDCs) have been attractive research topics due to their exotic electrical, optical, and mechanical properties and their suitability for usage in the future 2D electronic and optoelectronics [1–4]. This interest started after the first

extraction from graphite using the mechanical exfoliation method of graphene, and since then, graphene has become one of the most intensively studied topics in the field of 2D materials [5–8]. Although it attracts tremendous attention and it is used in electronic device development, graphene is an inconvenient substance for digital and optoelectronic applications due to the lack of optical bandgap [9]. This absence caused the focus of the 2D material research to shift to the TMDCs [10–13].

TMDCs are semiconductors in the form of  $\text{MX}_2$ , in which M is a transition metal atom and X is a chalcogen element. The metal atom in the middle bond covalently to the chalcogen atoms and this bonding constructs a sandwich form in the single layer [14]. Although they have a strong covalent bond in-plane, their out-of-plane bonding is weak due to the van der Waals interactions. Because of that they are able to stay stable, even in the monolayer form [15]. Also, with the help of this weak bonding in out-of-plane, they can be obtained as a monolayer by mechanical exfoliation called as the “Scotch-tape” method [16]. In addition, the synthesis of thin-film TMDCs has been achieved at wafer-scale by

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metal-organic chemical vapor deposition (MOCVD) [17], and small-scale with chemical vapor deposition (CVD) [18].

As a special TMDC semiconductor, monolayer molybdenum disulfide ( $\text{MoS}_2$ ) arouse the interest of many researchers [19–21]. Similar to other TMDCs, the  $\text{MoS}_2$  experiences an indirect to direct bandgap transition (from 1.3 to 1.8 eV) when it turns to monolayer from bulk [22]. Thus, this transition significantly improves photoluminescence and absorption dynamics of the material. Moreover, monolayer  $\text{MoS}_2$  with a thickness of 0.65 nm has a flexible mechanics due to its high Young's modulus [23], and it has also high carrier mobility and high on-off ratio up to  $10^9$  [24]. These properties promise flexible, high-performance, small, and light optoelectronic devices such as ultra-fast detector, field-effect transistors, and photovoltaic cells [25]. Especially high absorption coefficient and bandgap in the visible range makes  $\text{MoS}_2$  ideal for 2D optoelectronic device applications [19].  $\text{MoS}_2$  gives about 39- and 13-times higher absorption if it is compared with traditional materials such as Si and GaAs in 1-nm layer thickness. Despite its superiority over traditional materials in atomic scale dimension, monolayer  $\text{MoS}_2$  itself absorbs only  $\sim 10\%$  of incident light [26]. Weak absorption due to their extremely thin thickness is the major problem in the way of their use in photoconversion systems. Thus, this 2D semiconductor should be coupled with a trapping scheme to increase its light absorption [27].

To be able to enhance this absorption feature of the TMDCs towards the unity, three main trapping schemes can be employed: (i) metal-based plasmonic [28–33], (ii) guided mode resonance [34–37], and (iii) photonic crystal [38–40]-based designs. Although plasmonic structures are common solution to enhance the light-matter interaction in TMDC-based absorbers, metals cause parasitic absorption (Joule loss) due to their larger extinction coefficient (compared with that of TMDCs). Thus, the most of the light is consumed in metal parts and absorption efficiency within the TMDC monolayer is low. Moreover, they have large full-width-at-half-maximum (FWHM), due to the fact that the surface plasmon resonances (SPRs) have relatively broad spectral coverage [41]. Guided mode resonance-based designs can provide higher absorption efficiencies (due to their low-loss material nature); however, they have also large FWHM values. Photonic crystal-integrated designs have also bulky nature with multiple fabrication routes, which violates the 2D advantage of these TMDC based designs. Thus, in an ideal scheme, we need to have a lossless and 2D trapping architecture to maximize light coupling into TMDC monolayer. This can be achieved with all-dielectric-based metasurface designs that support the excitation of Mie resonances. Dielectric-based designs are lowloss alternatives of metallic designs and have found many applications, in recent years [42–46]. With the help of Mie resonance, the plasmonic effect can be replicated and the lossless characteristic

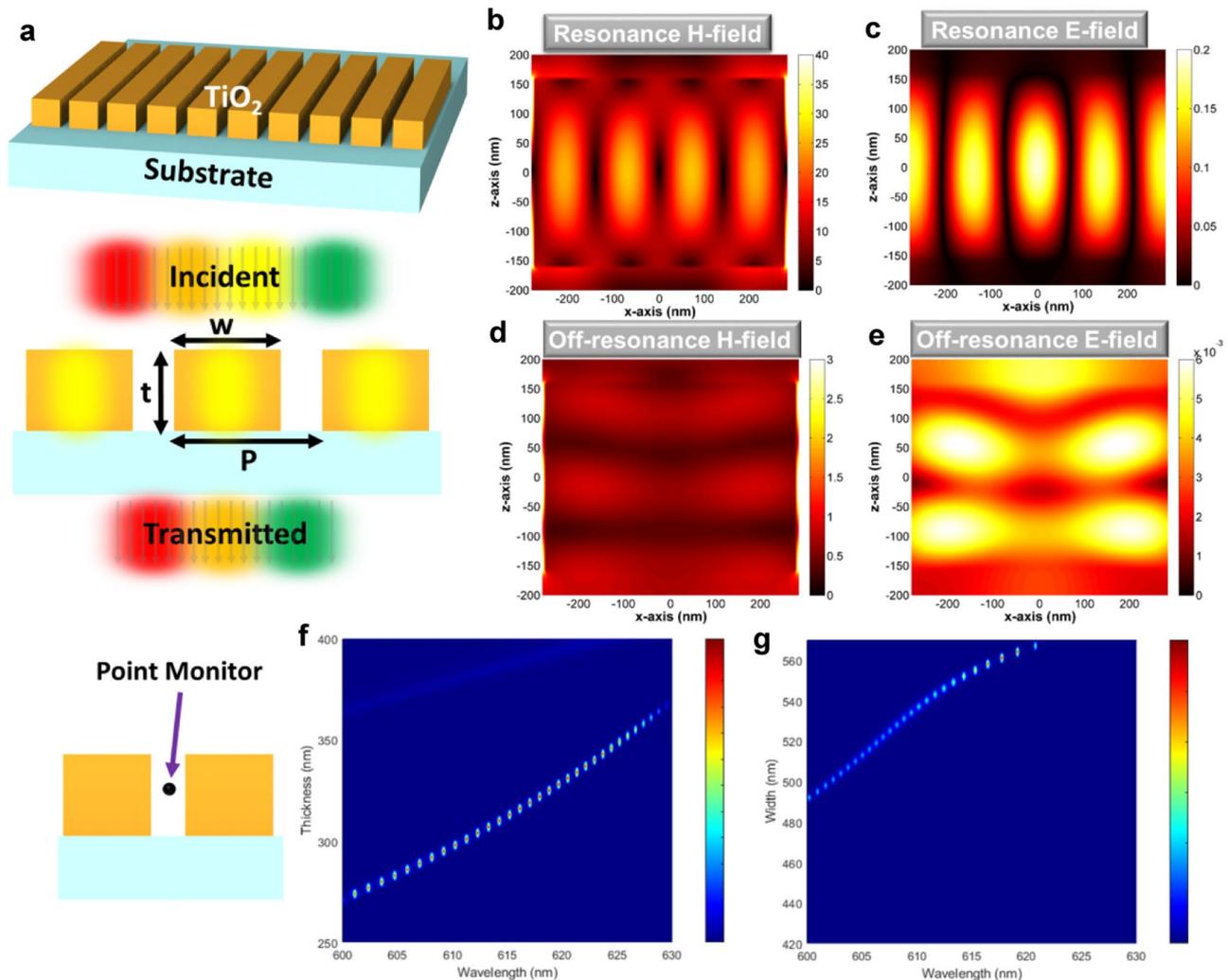
of the dielectrics can provide a narrow spectral range light confinement, with high absorption efficiency. In addition to these features, the strong light confinement in the dielectric medium can provide high sensitivity into refractive index change. This will provide the opportunity to manipulate the absorption peak response. In other words, this structure can provide thermal tunability without the use of any phase-change material.

In this paper, we propose a dielectric based metasurface platform to achieve ultra-narrowband light absorption within a monolayer thick TMDC layer. For this purpose, we first optimize the metasurface design itself. Then, we couple this design with mono and multilayer TMDCs to observe better absorption results. For this purpose, we choose  $\text{MoS}_2$ , and  $\text{WS}_2$  as the most commonly used TMDCs. The coupling of light into Mie resonances, supported by dielectric nanograting, provides narrowband absorption within the TMDC layer. To reach further enhancement, a cavity design is integrated into this dielectric-based metasurface. For the best optimized design, the absorptance efficiency reaches to 0.85 and FWHM stays as narrow as 3.1 nm. Finally, we move one step forward to show the thermal tunability characteristic of our design, without use of any phase change material. This is achieved due to strong light confinement within the design. Due to this confinement, any small change in the refractive index is seen by the resonant design. Thus, the resonance frequency shifts and thermal tunability is acquired. The thermal sensitivity of the above mentioned optimized design reaches to  $0.0096 \text{ nm}/^\circ\text{C}$ .

## Result and Discussion

Figure 1 a is a schematic illustration of the dielectric metasurface design, consisting of nano gratings from titanium dioxide ( $\text{TiO}_2$ ). Thickness and width of the  $\text{TiO}_2$ , and the periodicity of the structure are indicated as  $t$ ,  $w$ , and  $P$ , respectfully. The refractive index of  $\text{TiO}_2$  is set at 2.3705 at room temperature [47]. The calculations to find optimal geometries and dimensions are performed by a commercial finite-difference time-domain (FDTD) software package (Lumerical FDTD Solutions) [48]. While the numerical calculations are executed, the incident light propagation direction is selected as perpendicular to the  $x$ - $y$  plane and, boundary conditions for  $x$  and  $y$  directions are chosen as periodic boundary and perfectly matched layer (PML) for the  $z$  direction.

In standard circumstances,  $\text{TiO}_2$  cannot absorb or confine light due to its transparent and lossless nature in the visible range. Using the proposed design in Fig. 1a, Mie resonances can be excited, and it can provide a spectral selective strong confinement of the light. This strong confinement at the resonance frequency can be seen in

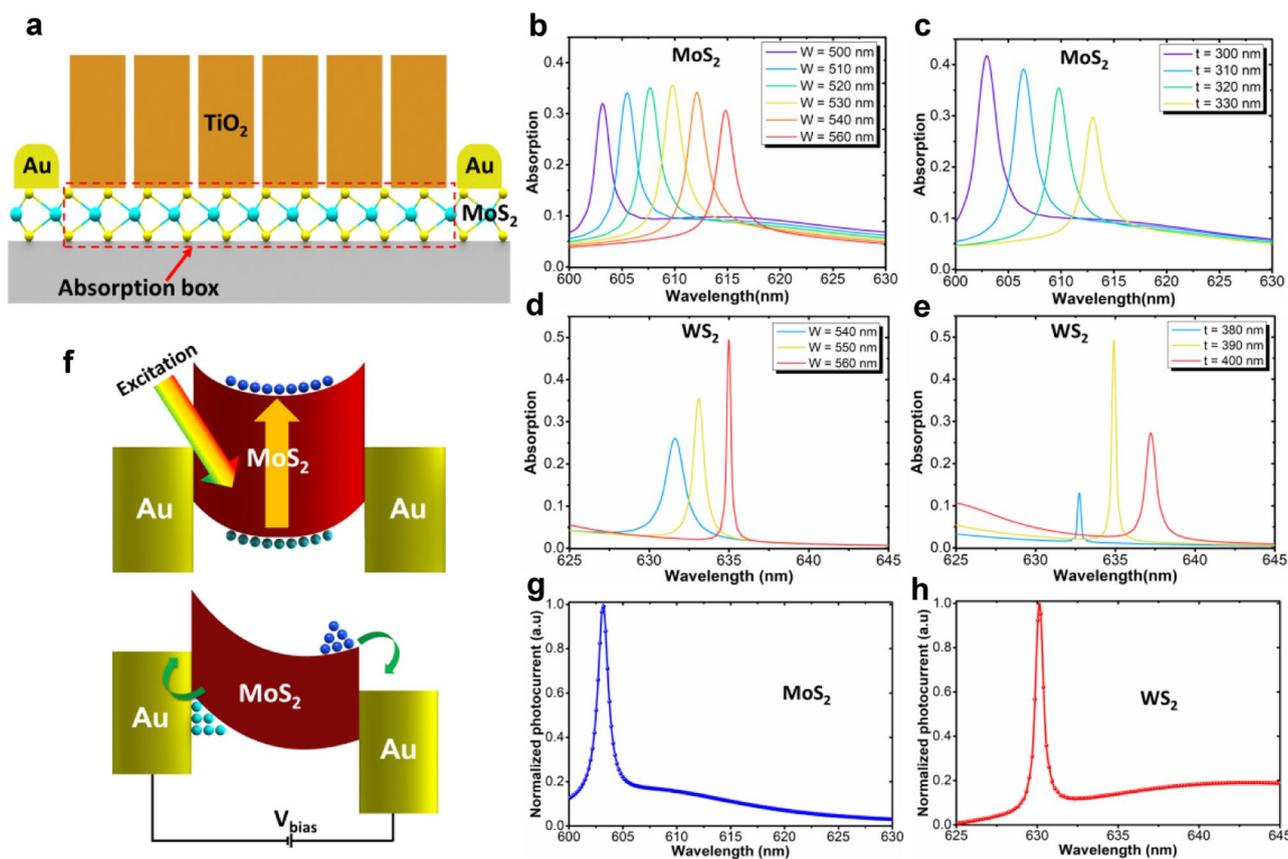


**Fig. 1** (a) Schematic representation of the proposed dielectric metasurface structure. H-field and E-field distributions across the design for (b, c) resonance frequency and (d, e) off-resonance region. Contour plots of magnetic field intensity dependence to (f) thickness and (g) width

the E-field and H-field profiles across the structure in the Fig. 1b, c. The structure becomes an ordinary transparent dielectric in the off-resonance frequencies, as shown in the Fig. 1d, e. The structure has an ability to provide field enhancement, at specific thickness and width values. For better qualitative comparison, the magnetic field intensities are probed at the point monitor position for the desired spectral range. These field intensities are studied for different  $t$  and  $W$  values, while  $P$  is kept as 570 nm in all simulations. The sweeping results are shown as a contour plot in the Fig. 1f, g. The contour plots imply that at specific geometrical values the structure provides large absorption cross section. Thus, if a material with high absorption coefficient is embedded into the bottom of the  $\text{TiO}_2$  grating, spectrally selective light absorption can be achieved, and the structure becomes an ultra-narrow semiconductor absorber.

Monolayer TMDCs are excellent candidates for this purpose. Due to their single atomic-scale thickness, the off-resonance absorption is quite low and this causes a high-contrast light absorption.  $\text{MoS}_2$  and  $\text{WS}_2$  as a monolayer from TMDCs are selected for the absorbent semiconductor layer. In their own case, these monolayer materials can only absorb small amount of the incoming power due to its ultrathin thickness of  $\sim 0.62$  nm. To be able to increase the absorbance magnitude, monolayers are coupled to this resonant metasurface design. This metasurface confines a specific wavelength and passes the rest of spectrum. In other words, it can be seen as an ultra-narrowband band-stop filter.

To examine the absorption features of the structure,  $\text{MoS}_2$  is placed right underneath the  $\text{TiO}_2$  nano-grating as a monolayer and absorption measurement is taken on the monolayer region. The proposed configuration is shown in the Fig. 2a. Starting with the optimal thickness and width values



**Fig. 2** (a) Schematic representation of the all-dielectric narrow-band absorber structure, absorption measurement for varying thickness and width of TiO<sub>2</sub> with MoS<sub>2</sub> (b–c) and WS<sub>2</sub> (d–e). (f) The band align-

ment between Au and MoS<sub>2</sub> layer before and after applying bias. The simulated photocurrent values generated by TMD monolayers in (g) MoS<sub>2</sub> and (h) WS<sub>2</sub> cases

of the metasurface from previous section, the dependence of spectral absorption within the TMDCs to geometrical values is analyzed with 10-nm intervals in Fig. 2b, c. For this purpose, the absorption within the TMDC layer is calculated using below Eq. (1).

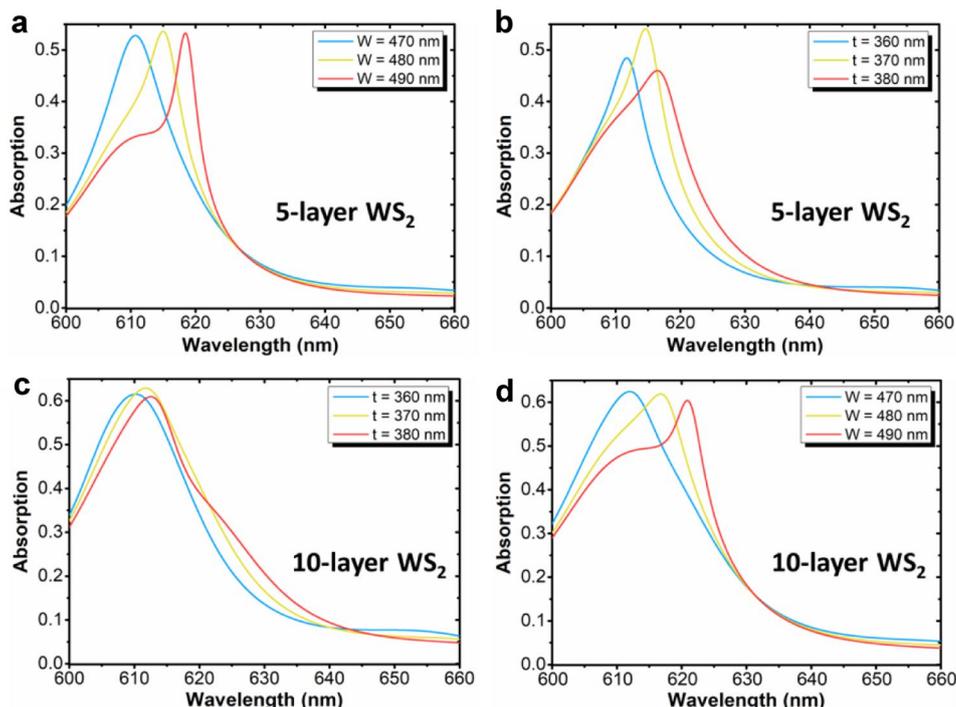
$$A(\lambda) = \frac{1}{P_{\text{incident}}} \iiint 0.5\omega \text{imag}(\epsilon_{\text{MoS}_2}) |E(x, y, z, \lambda)|^2 dx dy dz \tag{1}$$

This formula has been used to ensure that the absorption is fully within the layer and parasitic absorption is ignored. After finding the optimum dimensions for the width value, which is 530 nm, thickness optimization is performed with final result of 300 nm. For these geometries, the absorption resonance is located at 603 nm wavelength with a maximum absorptance 0.42 and FWHM of 2.7 nm. The same procedure is applied to the WS<sub>2</sub> monolayer. The strongest absorption peak is found at the 560 nm width and 390 nm thickness, as shown in the Fig. 2d, e. The spectral position, amplitude, and FWHM values of this design are 634.9 nm, 0.493, and 0.3 nm.

Thus, the function of proposed metasurface is the selective coupling of the incident light into the TMDC semiconductor monolayer. To convert this structure into the photodetector, Au contacts can be made on its surface to form metal-semiconductor-metal (MSM) configuration. In this design, upon broadband light illumination, mostly photons corresponding to resonance peak position are absorbed and a photocurrent spectrum similar to absorption profile is achieved. This phenomenon is illustrated in the Fig. 2f. COMSOL Multiphysics finite element method solver is used, to numerically analyze the photo-electrical response of the photodetector. As shown in Fig. 2g, h, the proposed design can provide an order of magnitude larger photocurrents in the resonance peak, compared with off-resonance ones. Thus, the proposed configuration can be utilized to accomplish an ultra-narrowband 2D photodetector.

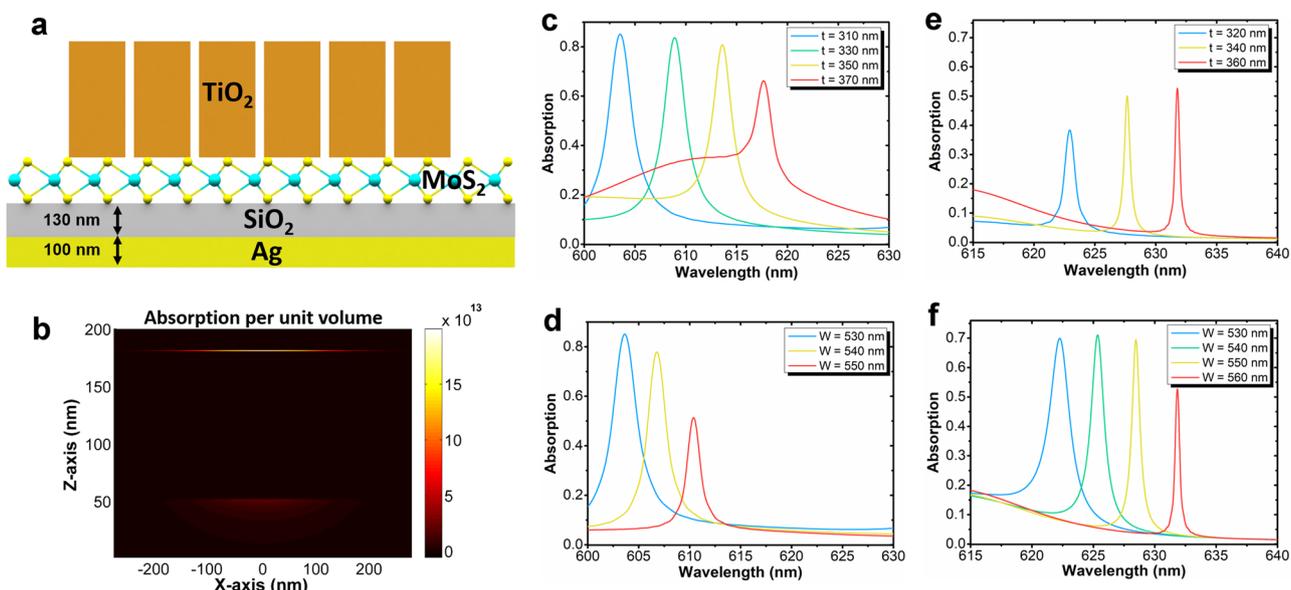
While the absorption peak for both MoS<sub>2</sub> and WS<sub>2</sub> is quite narrow FWHM, the absorptance amplitude is much below unity. To be able to increase the amplitude towards the unity, two approach can be proposed: (i) increasing the semiconductor layer thickness and (ii) using cavity design to improve light trapping. To examine the first idea, 5-layer, and 10-layer WS<sub>2</sub>

**Fig. 3** Absorption measurement for varying (a) thickness and (b) width of TiO<sub>2</sub> with 5 Layer WS<sub>2</sub> and varying (c) thickness and (d) width of TiO<sub>2</sub> with 10 Layer WS<sub>2</sub>

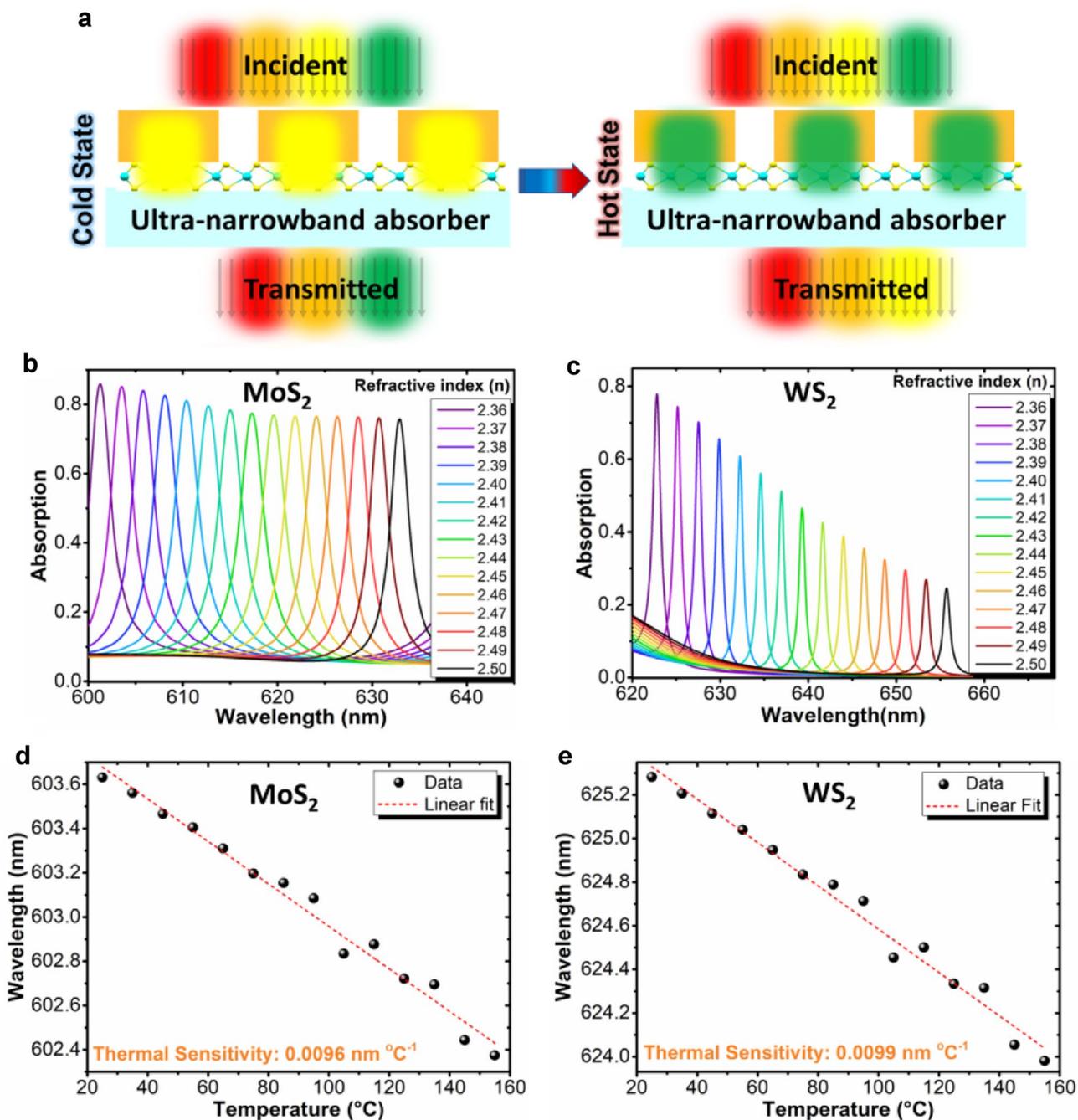


are placed instead of the monolayer. The similar optimization process is conducted to get the highest absorption peak. The results are presented in Fig. 3a, b, c, d. The maximum value is increased up to 0.55 with 5-layer and 0.63 with 10-layer configurations. However, this improvement is achieved in the expense of broadening of FWHM (15.1 nm for 5-layer, and 23.1 nm for 10-layer), and increasing the background absorption, compared

with that of monolayer (see Fig. 2d, e). This higher off-resonance absorption is the result of thicker semiconductor layer. Thus, in this configuration, although the amplitude has been gradually improved, the selectivity (the ratio of resonance and off-resonance absorptions) is decreased and FWHM is wider. For the second approach, the proposed metasurface design is placed on top of a metal-insulator (MI) cavity. Fig. 4(a) is an



**Fig. 4** (a) Proposed optimized cavity design and (b) its absorption profile across the cavity with MoS<sub>2</sub> and absorption measurement for varying thickness and width of TiO<sub>2</sub> with MoS<sub>2</sub> (c–d) and (e–f) WS<sub>2</sub>



**Fig. 5** (a) Schematic illustration of red-shift of the absorption peak as the temperature is increased. Absorption measurement for varying refractive index of TiO<sub>2</sub> with (b) MoS<sub>2</sub> and (c) WS<sub>2</sub> in the proposed

design and thermal sensitivity plots with different temperature for (d) MoS<sub>2</sub> and (e) WS<sub>2</sub>

illustration of the cavity enhanced design. Silver (Ag) is chosen as the bottom metal mirror. As shown in the absorption per unit volume plot, the dominant part of the absorption occurs within the TMDC monolayer and minor part is lost as parasitic absorption (Jule loss) within the bulk metal. The width and thickness value of TiO<sub>2</sub> is optimized in Fig. 4(c) and Fig. 4(d) for MoS<sub>2</sub>

and Fig. 4(e) and Fig. 4(f) for WS<sub>2</sub>. The optimal absorption values are much higher compared to previous designs. For MoS<sub>2</sub> case, cavity design has improved absorbance to a value as high as 0.85 with a FWHM of 3.1 nm. While, WS<sub>2</sub> cavity design absorbance value is reached to 0.72 at with a FWHM as narrow as 1.4 nm.

Besides having high magnetic field enhancement, the structure is also sensitive to the environment temperature. As schematically illustrated in Fig. 5a, the structure has thermally tunable absorption response. This is due to strong field confinement within the nanogratings. Changing temperature cause the change in the refractive index of the  $\text{TiO}_2$  material. Although  $\text{TiO}_2$  has very small thermo-optic coefficient of  $\sim 4.2 \times 10^{-5}$ , the formation of hot spots in small volume fractions triggers the high sensitivity of the resonance mode into environmental index change. Thus, the change in temperature is resulted with the shift in the absorption spectrum. To show this effect refractive index of  $\text{TiO}_2$  is swept from 2.36 to 2.5 with 0.01 intervals. This effect is shown with both  $\text{MoS}_2$  and  $\text{WS}_2$  material in the Fig. 5a, b. These results are indicated that the design can act as an ultra-narrowband single-color detector, without using any phase change material. The thermal sensitivity graphics of the cavity design for each monolayer are also shown in the Fig. 5d, e. The sensitivity values for  $\text{MoS}_2$  and  $\text{WS}_2$ -based designs are  $0.0096 \text{ nm}^\circ\text{C}$  and  $0.0099 \text{ nm}^\circ\text{C}$ , respectively. These values are significantly high compared with inherent small thermo-optic coefficient of  $\text{TiO}_2$ .

## Conclusion

This study is focused to enhance the absorption of TMDCs specifically  $\text{MoS}_2$  and  $\text{WS}_2$  monolayer with the help of the all-dielectric metasurface structure. The absorption behavior of the proposed structure is an ultra-narrowband response over the visible range. The cavity design and multilayer approaches are studied for the further enhancement of the absorptance amplitude. The cavity approach was shown to be a promising approach to maximize the performance of the design. Due to strong field confinement, the proposed design revealed thermally tunable absorption response, without the use of phase-change component. This paper proposes a universal design approach that can be used in other 2D and ultrathin optical systems where strong and selective light absorption can be achieved in atomic scale dimensions. The tunability added to this design provides the opportunity to realize multi-color photodetectors using the same design. As temperature changes, the absorption peak experiences a spectral shift and by this way we can detect different colors using the same photodetector design. Besides, this design has the ability to be used in sensing application as well. As an unknown agent is coated on the grating surface, the resonance peak will shift, and from the produced shift value, one can estimate the refractive index of the material.

**Authors' Contributions** First author (E.B.) carried out design, simulation, and analysis of the results. Second author (A.G.) helped in simulation and discussion steps of the paper. Last author (E.O.) supervised the study. All the authors contributed in the paper writing.

## Compliance with Ethical Standards

**Competing Interests** The authors declare that they have no competing interests.

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