



PAPER

Investigation of angstrom-thick aluminium oxide passivation layers to improve the gate lag performance of GaN HEMTs

RECEIVED
30 April 2019REVISED
21 June 2019ACCEPTED FOR PUBLICATION
4 July 2019PUBLISHED
12 July 2019Melisa Ekin Gulseren^{1,2} , Gokhan Kurt² , Turkan Gamze Ulusoy Ghobadi^{3,4}, Amir Ghobadi^{1,2}, Gurur Salkim², Mustafa Ozturk², Bayram Butun² and Ekmel Ozbay^{1,2,3,5}¹ Department of Electrical and Electronics Engineering, Bilkent University, 06800 Ankara, Turkey² Nanotechnology Research Center, Bilkent University, 06800 Ankara, Turkey³ UNAM-Institute of Materials Science and Nanotechnology, Bilkent University, 06800 Ankara, Turkey⁴ Department of Energy Engineering, Faculty of Engineering, Ankara University, 06830 Ankara, Turkey⁵ Department of Physics, Bilkent University, 06800 Ankara, TurkeyE-mail: melisa.gulseren@bilkent.edu.tr**Keywords:** AlGaN, atomic layer deposition, dielectric, GaN, gate lag, HEMT, passivation**Abstract**

In this paper, we report an angstrom-thick atomic layer deposited (ALD) aluminum oxide (Al_2O_3) dielectric passivation layer for an AlGaN/GaN high electron mobility transistor (HEMT). Our results show a 55% improvement in the gate lag performance of the design and a decrease by half in interface state density upon coating with two cycles of ALD Al_2O_3 . DC characteristics such as current density, threshold voltage, and leakage currents were maintained. ALD Al_2O_3 passivation layers with thicknesses up to 10 nm were investigated. XPS analyses reveal that the first ALD cycles are sufficient to passivate GaN surface traps. This study demonstrates that efficient passivation can be achieved in atomic-scale with dimensions much thinner than commonly used bulk layers.

1. Introduction

GaN based high electron mobility transistors (HEMTs) are advantageous for high power and high frequency applications due to their superior properties such as wide band gap, high breakdown field, and high saturation velocity [1–3]. Although GaN HEMTs are widely preferred for a variety of commercial and military applications, unpassivated devices suffer from limitations in characteristics such as radio frequency (RF) current, transconductance, and breakdown voltage as a result of the trapping of electrons in the surface states created by defects and dislocations [4, 5]. Surface passivation has become the most established approach to address the problem of current collapse, due to its simplicity and efficacy, with SiN_x being the most commonly preferred dielectric [6–8]. SiN_x passivation schemes are particularly effective for III-N devices due to the relatively low state density at the SiN_x /III-N interface and the possibility of *in situ* MOCVD deposition [9]. Numerous other dielectrics have been demonstrated as passivation layers over the years with different advantages [10–12]. Al_2O_3 passivation has been demonstrated with record drain current and transconductance and improved pulsed I-V characteristics [13–15]. Al_2O_3 is an attractive material because of its large bandgap, relatively high dielectric constant, and high breakdown field and has the advantage of high quality oxide/III-N interfaces when atomic layer deposited (ALD) [16].

Recently, angstrom-thick passivation has been demonstrated as an efficient route to provide both surface coverage and trap passivation in various applications [17–20]. It was found that the first ALD cycles will only passivate surface defects without introducing any new bulk traps. Thus, it is envisioned that this atomic-scale selective passivation is the optimal way for GaN HEMTs.

In this paper, we demonstrate that two ALD cycles of Al_2O_3 will lead to about 55% improvement in the gate lag characteristic of the GaN HEMT compared to that of an unpassivated design. This improvement has been correlated with the reduction in the density of GaN surface traps, which has been studied with frequency dependent conductance measurements. An in-depth XPS surface analysis confirmed that 2 Å Al_2O_3 is sufficient

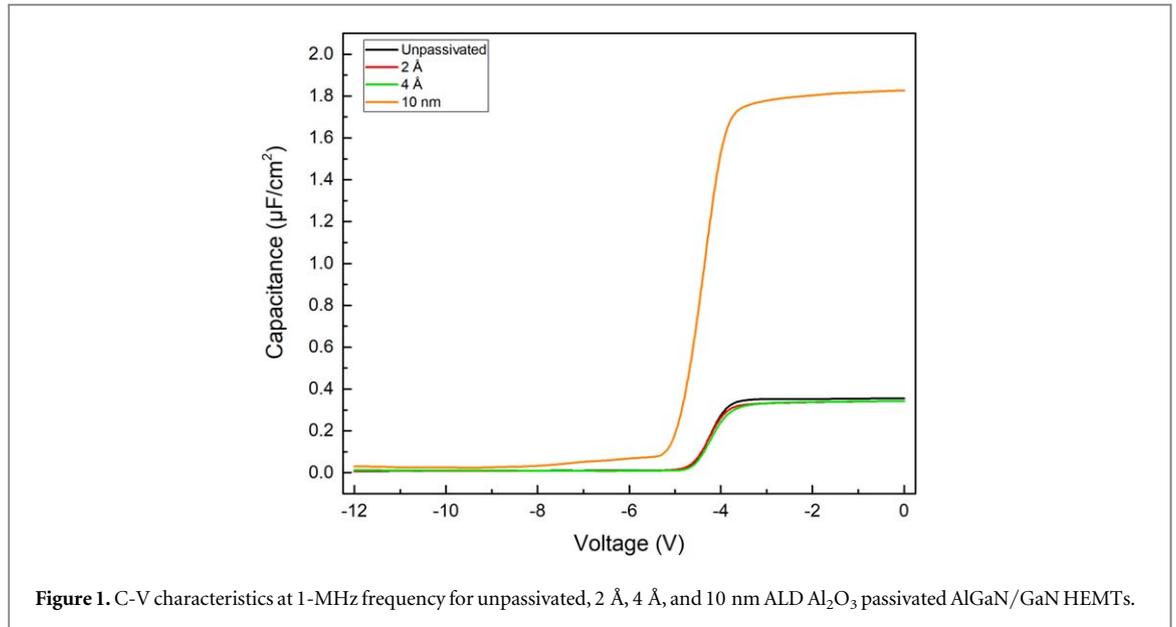


Figure 1. C-V characteristics at 1-MHz frequency for unpassivated, 2 Å, 4 Å, and 10 nm ALD Al₂O₃ passivated AlGaIn/GaN HEMTs.

for passivation of the AlGaIn surface. DC characteristics such as knee voltage and transconductance were improved and other characteristics were maintained for the passivated device.

2. Device fabrication

The epitaxial structure consists of a 20 nm AlN nucleation layer and 1350 nm Carbon doped high resistive buffer layer on a semi-insulating SiC wafer, followed by 150 nm high mobility GaN channel, 1 nm AlN spacer, and 22 nm AlGaIn barrier layer with 26% Aluminum content capped with 3 nm GaN. The electron mobility and 2DEG density were found to be 2061 cm² V⁻¹.s and 1.07 × 10¹³ cm⁻² respectively. Average sheet resistivity was 284 Ohms/Sq across the 3 inch wafer. Device fabrication began with mesa isolation. Ohmic contacts were formed with electron beam deposited Ti/Al/Ni/Au stack, followed by 30 s rapid thermal annealing at 850 °C in N₂ ambient. 250 nm gate contacts were formed with electron beam lithography and Ni/Au gate metal deposition. For device passivation Al₂O₃ depositions were carried out at 200 °C in an ALD reactor (Cambridge Nanotech Savannah S100) employing Al(CH₃)₃ solution as the deposition precursor. The pulse and purge durations were 0.015 and 10 s, respectively. Water was used as the oxygen precursor. The deposition rate was found as 1 Å/cycle. Different devices were fabricated with 2 cycles, 4 cycles, 10 cycles, and 100 cycles to obtain passivation layer thicknesses of 2 Å, 4 Å, 1 nm, and 10 nm, respectively. The devices in this letter have a gate width of 2 × 125 µm, a gate-source spacing of 1.5 µm, and gate drain spacing of 3.25 µm. Schottky diodes were fabricated with a similar process flow for capacitance measurements.

3. Results and discussion

Capacitance-voltage (C-V) measurements of fabricated Schottky diodes were conducted in order to examine the gate characteristics of the different passivation schemes. Figure 1 shows the C-V characteristics for the unpassivated, 2 Å, 4 Å, and 10 nm ALD Al₂O₃ passivated samples measured at a frequency of 1 MHz. For the unpassivated, 2 Å, and 4 Å passivated samples similar capacitance curves are obtained. An increase in capacitance is observed for the 10 nm Al₂O₃ passivated sample, confirming the formation of the oxide layer. A negative threshold shift is also observed for the sample with 10 nm passivation.

Frequency dependent conductance measurements were carried out to determine the interface state density D_{it} . The frequency was varied from 10 kHz to 2 MHz. The parallel C_m — G_m circuit was considered for the equivalent circuit to the interface. The interface state density was obtained from [21] as

$$\frac{G_p}{\omega} = \frac{q\omega\tau_{it}D_{it}}{1 + (\omega\tau_{it})^2} = \frac{\omega G_m C_{ox}^2}{G_m^2 + \omega^2(C_{ox} - C_m)^2} \quad (1)$$

where C_m and G_m are the measured capacitance and conductance at frequency ω , and C_{ox} is the oxide capacitance. The calculated G_p/ω values for the unpassivated, 2 Å, 4 Å, and 10 nm ALD Al₂O₃ passivated schemes are shown in Figure 2. The extracted D_{it} values are given in Table 1. The relative position of the trap energy level

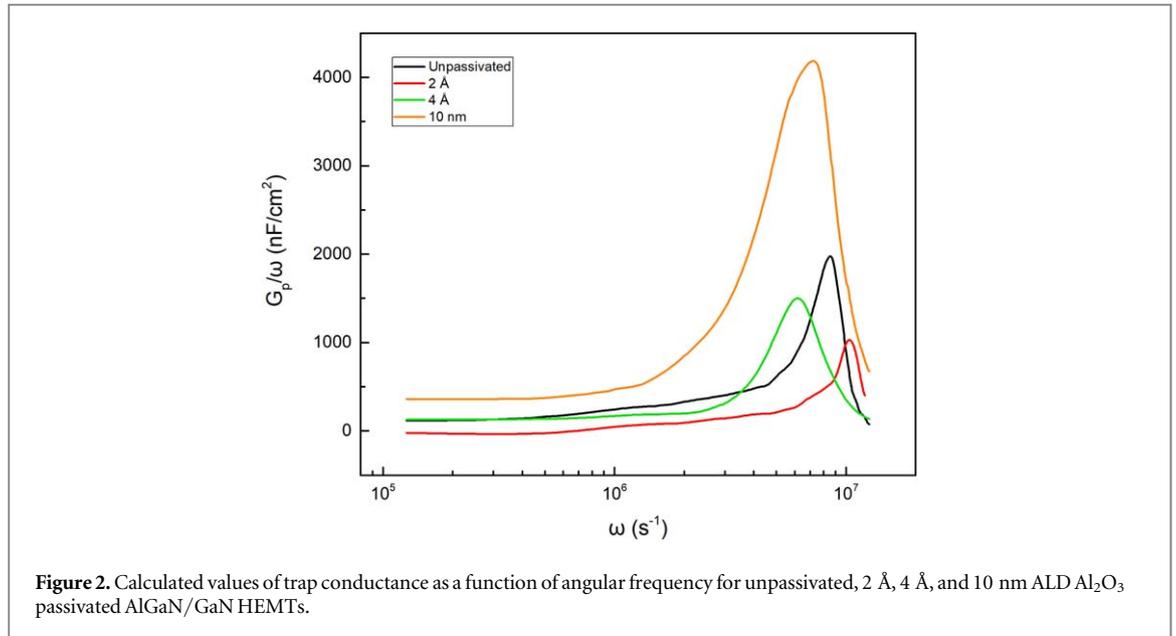


Figure 2. Calculated values of trap conductance as a function of angular frequency for unpassivated, 2 Å, 4 Å, and 10 nm ALD Al₂O₃ passivated AlGaIn/GaN HEMTs.

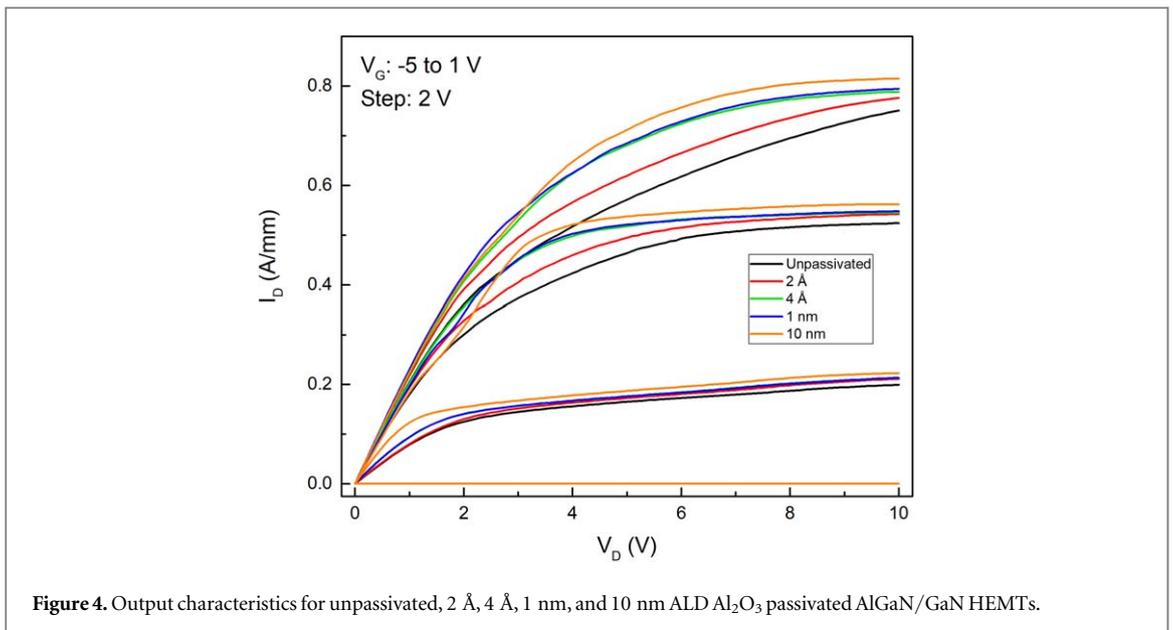
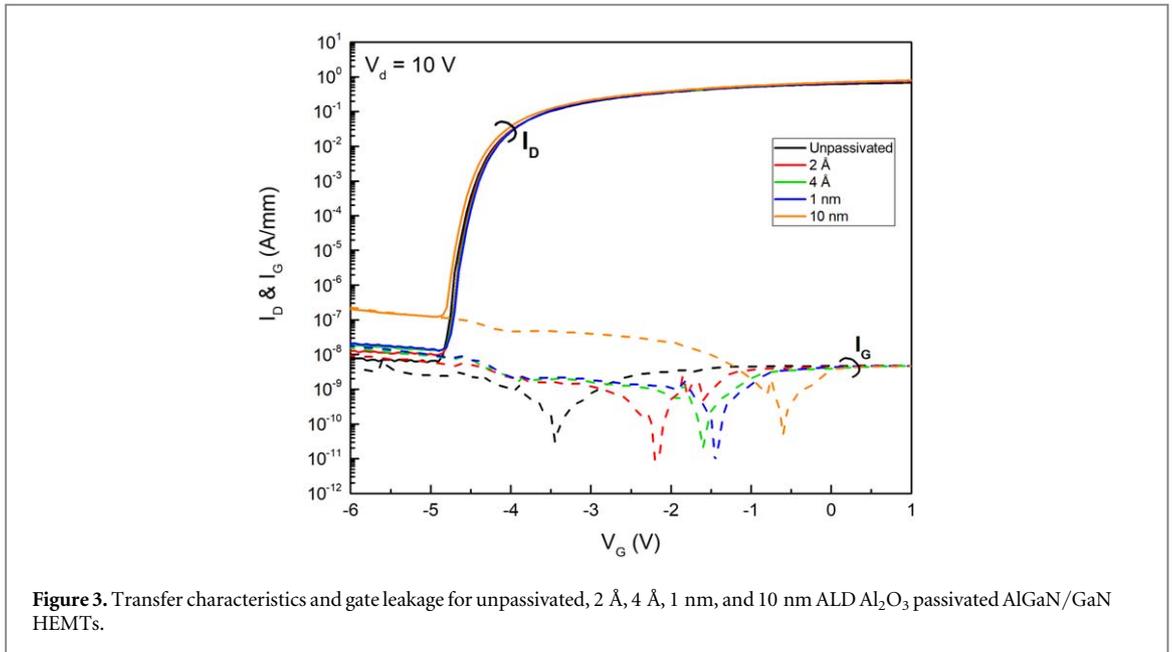
Table 1. List of measurements for devices with no passivation, 2 Å, 4 Å, 1 nm, and 10 nm ALD grown Al₂O₃.

Parameter	Units	Unpassivated	2 Å	4 Å	1 nm	10 nm
D_{it}	$\text{eV}^{-1}\text{cm}^{-2}$	3.08×10^{13}	1.62×10^{13}	2.36×10^{13}	—	6.52×10^{13}
I_{dss}	A mm^{-1}	0.75	0.78	0.79	0.79	0.82
V_{knee}	V	8.6	8	6.8	6.8	6.5
V_{th}	V	-4.06	-4.07	-4.07	-4.07	-4.15
g_m	mS mm^{-1}	179	186	187	188	188
$I_{g,leak}$	nA mm^{-1}	2.7	3.4	1.9	1.6	21.4
$I_{d,leak}$	nA mm^{-1}	8.3	12.7	19.3	20.3	206
Gate Lag	%	20.3	9.4	11.0	11.8	8.5

with respect to the conduction band edge is found to be in the range of 0.18–0.19 eV and the time constants in the range of 0.19–0.32 μs for the samples. An interface state density of $3.08 \times 10^{13} \text{ eV}^{-1} \text{ cm}^{-2}$ is obtained for the unpassivated sample. For the 2 Å Al₂O₃ sample the lowest D_{it} of the studied samples is obtained, with a value of $1.62 \times 10^{13} \text{ eV}^{-1} \text{ cm}^{-2}$, which corresponds to nearly a 50% decrease in D_{it} compared to the unpassivated sample. For the 4 Å Al₂O₃ sample a slight increase is observed in the interface state density, however the obtained value of $2.36 \times 10^{13} \text{ eV}^{-1} \text{ cm}^{-2}$ is still lower than that achieved for the unpassivated sample. For the 10 nm Al₂O₃ sample, however, the highest D_{it} of $6.52 \times 10^{13} \text{ eV}^{-1} \text{ cm}^{-2}$ is calculated, which corresponds to a nearly 50% increase compared to the unpassivated sample. Thus, the conductance method measurements indicate that angstrom thick ALD passivation layers are able to passivate surface states selectively whereas thicker layers lead to increased surface state density.

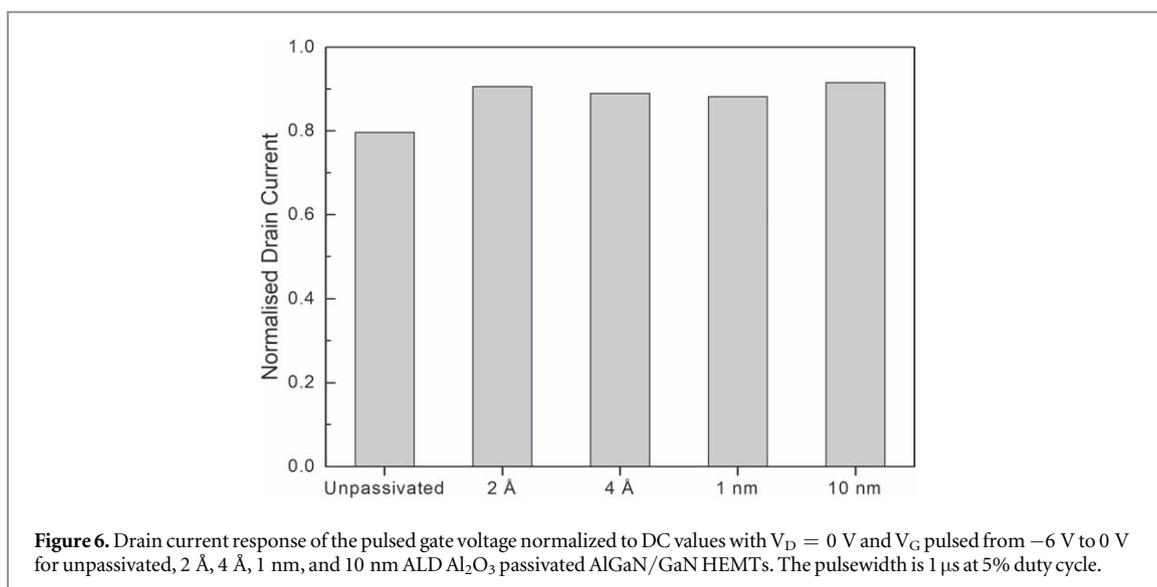
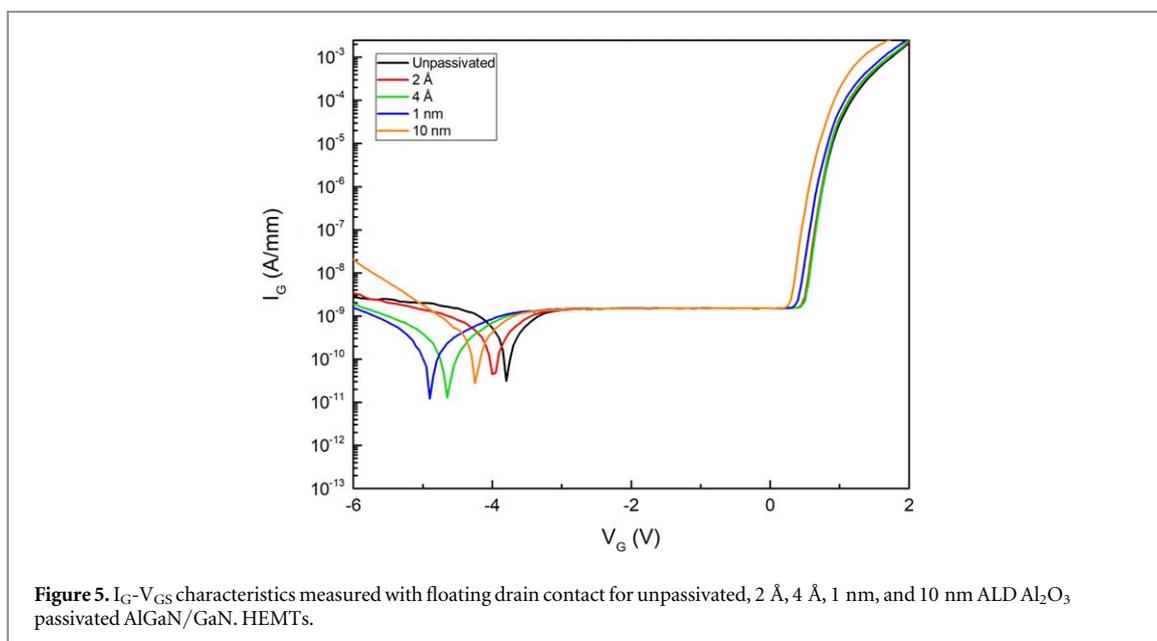
Table 1 lists the DC measurement results and calculated gate lag percentages for different devices, that are, unpassivated, 2 Å, 4 Å, 1 nm, and 10 nm ALD deposited Al₂O₃. Figure 3 compares the transfer characteristics and gate currents of the devices for each passivation scheme, at $V_D = 10 \text{ V}$. The threshold voltage was obtained as -4.06 V for the unpassivated device using linear extrapolation at maximum linear slope and does not vary significantly for 2 Å, 4 Å, and 1 nm Al₂O₃ passivation. For 10 nm ALD Al₂O₃ passivation, V_{th} decreases to -4.15 V, analogous to the decrease observed in the C-V measurement, due to the increased sheet charge from the passivation layer. The measured output characteristics are shown in Figure 4. The drain current density I_{dss} (at a gate voltage of $V_G = 1 \text{ V}$) was obtained as 0.75 A mm^{-1} for the unpassivated reference device, which increased slightly by 5% to a value of 0.79 A mm^{-1} for the 4 Å and 1 nm passivation scheme and by 10% to a maximum value of 0.82 A mm^{-1} for the 10 nm passivation scheme. Similarly, the peak transconductance value increased by 4.7% from 179 mS mm^{-1} for the reference device to 188 mS mm^{-1} for the devices with 1 nm and 10 nm passivation. An improvement in knee voltage is observed with increasing Al₂O₃ passivation thickness. The reason for these improvements in the drain current characteristics is related to enhanced electron transport caused by an increase in sheet carrier concentration, due to the reduction in surface states.

The drain leakage current $I_{d,leak}$ is extracted from the transfer characteristics measurements as the current at $V_G = -6 \text{ V}$ and $V_D = 10 \text{ V}$. $I_{d,leak}$ increases monotonically with passivation layer thickness, which is also



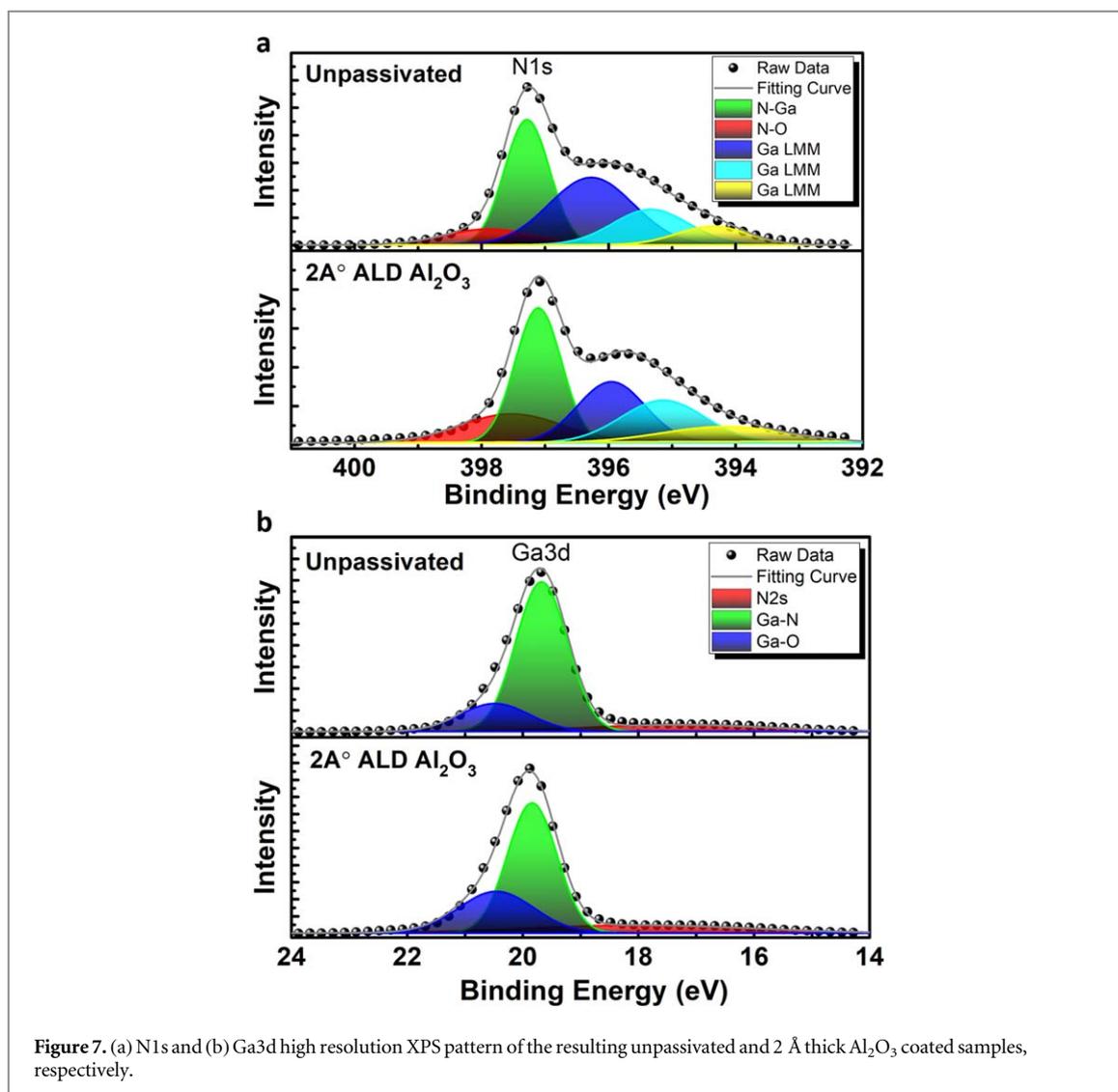
attributed to the increasing sheet charge of the 2DEG. The Schottky characteristics of the AlGaN/GaN HEMTs were measured with a floating drain contact (Figure 5) and the gate leakage current $I_{g,leak}$ is obtained from this measurement as the gate current at $V_G = -6$ V. The reverse gate leakage current, in contrast to the drain leakage current, does not follow a monotonically increasing trend, indicating that the reverse gate leakage current is not dominated by surface leakage. The reverse gate leakage current increases for the initial passivation 2 Å, decreases for the two subsequent passivation thicknesses, then increases to its highest value for the greatest passivation thickness of 10 nm. The dominant mechanism of gate leakage is attributed to be edge gate leakage current, as the trend of the reverse gate leakage current for increasing passivation thickness correlates inversely with the trend in the gate lag characteristic. A decrease in gate lag corresponds to an increase in gate leakage, indicating that extension of the gate depletion region towards the drain by the virtual gate effect acting to disperse the electric field lines induced by the depletion region reduces the edge current [22]. Passivation that suppresses the virtual gate thereby leads to an increase in gate leakage current. Low forward gate leakage current and nearly constant gate turn on voltage is maintained for each passivation scheme.

To further understand the correlation between the passivation layer and the surface states gate lag measurements were carried out. The gate lag measurements were performed using an Agilent E3631A power supply, Keysight Technologies 33500B waveform generator, and a Keysight InfiniVision DSOX2004A

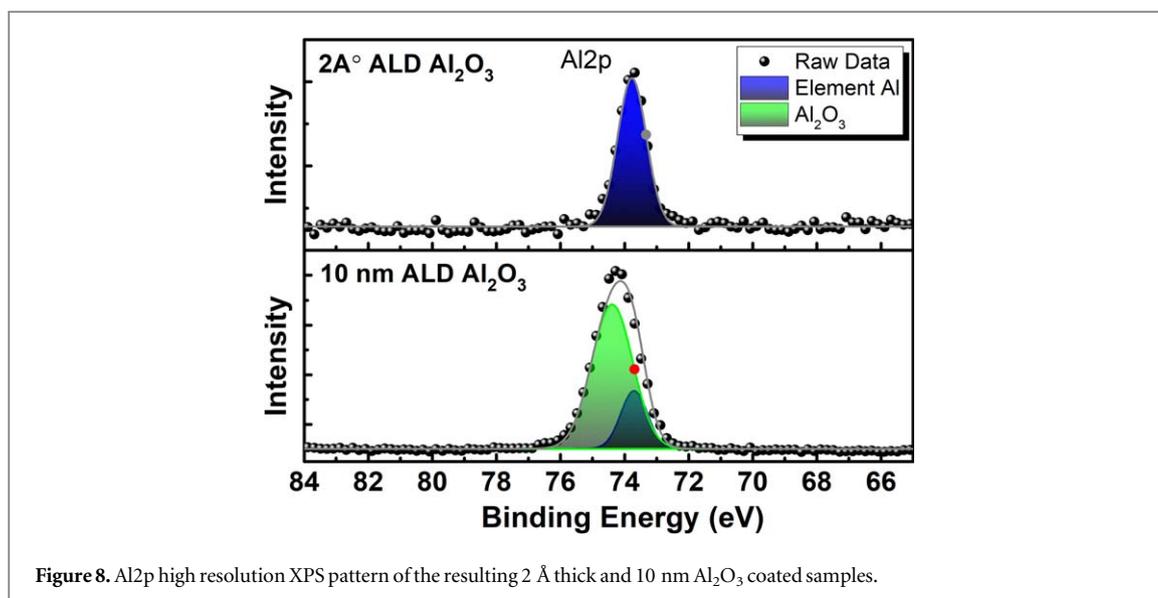


oscilloscope. The measurements were carried out by pulsing the gate from -6 V to 0 V, with a pulse width of $1 \mu\text{s}$ and period of $20 \mu\text{s}$ with the drain kept at 0 V. The pulsed drain current response with respect to the DC values are compared in Figure 6 for the studied passivation schemes. A gate lag of 20.9% is obtained for the unpassivated device, which improves by 55% to a gate lag value of 9.5% for the device with 2 Å passivation. For the 4 Å and 1 nm Al_2O_3 passivation layers, a slight increase in the gate lag characteristics is observed, however, significant improvement compared to the unpassivated case is maintained. For 10 nm Al_2O_3 passivation layer the lowest gate lag of 8.7% is observed. [23] reports that for ALD Al_2O_3 passivation thicknesses between 10 nm and 40 nm, increasing passivation thickness leads to decreasing current collapse, enabling the reduction of the current collapse below 15%. The results of this study indicate the same trend if passivation layers greater than 10 nm were studied, however, it is demonstrated that similar reduction in current collapse can be achieved with passivation layers as thin as 2 Å.

To gain insight on the passivation property of ALD Al_2O_3 layer, x-ray photoelectron spectroscopy (XPS) is utilized for unpassivated and 2 Å Al_2O_3 coated samples. Before ALD process, all the samples have been treated with a diluted HCl acid to remove the natural oxide layer. Figure 7 shows the N1s and Ga3d spectra of these two samples. Peak position correction was calibrated by referencing the C1s orbital peak position (284.8 eV) and the other peaks in the spectrum were shifted accordingly. As previously explained in several reports, the dominant surface defects of GaN are Ga and N vacancies (or dangling bonds). Based on the calculation of free energy by classical nucleation theory, most of the oxygen-derived hydroxyl groups such as OH radicals and H_2O or O_2 will



be chemisorbed near imperfections such as dangling bonds and vacancies. Looking back into Figure 7(a), the N1s spectra is deconvoluted into five main peaks [24]; three of which are assigned into Auger Ga LMM peaks, the dominant one comes from the Ga-N bond, and the one in the higher energy tail is attributed to N-O bonds. As we can see, the passivated sample has a stronger N-O related peak. During the ALD process, the Ga vacancy positions are passivated by oxygen molecules and form these N-O bonds. On the other side, the Ga3d spectra can be scrutinized to gain a further insight on the surface properties of the GaN layer. As illustrated in Figure 7(b), this spectrum is deconvoluted into three Gaussian profiles [25]; a broad and weak response originated from N2s orbitals, a dominant peak assigned to Ga-N bond, and a high energy response from Ga-O bonds. Looking at these peaks, it can be clearly seen that the Ga-O related peak is more dominant for 2 Å Al₂O₃ passivated case (compared to the unpassivated one). Therefore, similar to N1s data, the Al₂O₃ passivation layer will substitute N vacant positions and facilitate formation of Ga-O bonds. All the above-mentioned results confirm the efficient passivation of surface traps upon coating with the 2 Å Al₂O₃ layer. Oxygen-containing gas molecules tend to be chemisorbed on the surface of a semiconductor host through the capture of free electrons. Consequently, these chemisorbed radicals reduce the density of free carriers in the vicinity of the semiconductor surface and deplete the surface electron states. This, in turn, triggers the existence of the space charge region and induces band bending near the interface. On the other hand, these chemisorbed oxygen molecules are likely attached into trap states and dangling bonds. The passivation of surface traps, in first ALD cycles, reduces the surface traps density, mitigates adsorption of oxygen radicals, and consequently reduces band bending. However, as we go to larger ALD cycles, bulk trap states in Al₂O₃ layer start to become dominant and therefore diminish the abovementioned characteristics. Moreover, the Al2p spectra have been extracted for 2 Å and 10 nm coated samples. As Figure 8 clearly illustrates, the portion related to Al element is the major peak for the thin passivated sample. However, as we increase the Al₂O₃ layer thickness to 10 nm, the spectrum is mainly attributed to the



oxide related peak. Thus, the first two ALD cycle just passivates the defect states and does not form a continuous Al₂O₃ layer. However, the subsequent cycles trigger the formation of a continuous layer.

4. Conclusion

In conclusion, an angstrom-thick ALD Al₂O₃ dielectric is reported for the passivation of surface traps in AlGaIn/GaN HEMTs and passivation thicknesses up to 10 nm are studied. For the a passivation layer as thin as 2 Å Al₂O₃, a greater than 50% improvement in gate lag compared to the unpassivated device is achieved. Improvements in DC characteristics such as drain current density, knee voltage, and transconductance are observed. XPS analysis confirms that Ga vacancy positions are are passivated by oxygen molecules for Al₂O₃ as thin as 2 Å.

Acknowledgments

This work is supported by the TUBITAK under Project No. 116F041. One of the authors (E.O.) also acknowledges partial support from the Turkish Academy of Sciences.

ORCID iDs

Melisa Ekin Gulseren  <https://orcid.org/0000-0002-9375-2337>

Gokhan Kurt  <https://orcid.org/0000-0002-1210-4013>

Bayram Butun  <https://orcid.org/0000-0003-0892-4681>

References

- [1] Mishra U K, Parikh P and Wu Y-F AlGaIn/GaN HEMTs-an overview of device operation and applications *Proc. IEEE* **90** 1022–31
- [2] Chen K J *et al* GaN-on-Si power technology: devices and applications *IEEE Trans. Electron Devices* **64** 779–95
- [3] Chung J W, Hoke W E, Chumbes E M and Palacios T AlGaIn/GaN HEMT with 300-GHz f_{max} *IEEE Electron Device Lett.* **31** 195–7
- [4] Lu W, Kumar V, Schwindt R, Piner E and Adesida I 2002 A comparative study of surface passivation on AlGaIn/GaN HEMTs *Solid-State Electronics* **46** 1441–4
- [5] Vertiatchikh A V and Eastman L F 2003 Effect of the surface and barrier defects on the AlGaIn/GaN HEMT low-frequency noise performance *IEEE Electron Device Lett.* **24** 535–7
- [6] Edwards A P, Mittereder J A, Binari S C, Katzer D S, Storm D F and Roussos J A 2005 Improved reliability of AlGaIn-GaN HEMTs using an NH₃ plasma treatment prior to SiN passivation *IEEE Electron Device Lett.* **26** 225–7
- [7] Lee J-S, Vescan A, Wieszt A, Dietrich R, Leier H and Kwon Y-S Small signal and power measurements of AlGaIn/GaN HEMT with SiN passivation *Electron. Lett* **37** 130–2
- [8] Green B M, Chu K K, Chumbes E M, Smart J A, Shealy J R and Eastman L F The effect of surface passivation on the microwave characteristics of undoped AlGaIn/GaN HEMTs *IEEE Electron Device Lett.* **21** 268–70
- [9] Yatabe Z, Asubar J T and Hashizume T 2016 Insulated gate and surface passivation structures for GaN-based power transistors *J. Phys. D: Appl. Phys.* **49** 393001

- [10] Liu C, Chor E F and Tan L S 2007 Enhanced device performance of AlGaIn/GaN HEMTs using HfO₂ high-k dielectric for surface passivation and gate oxide *Semicond. Sci. Technol.* **22** 522
- [11] Kordoš P, Kúdela P, Gregušová D and Donoval D 2006 The effect of passivation on the performance of AlGaIn/GaN heterostructure field-effect transistors *Semicond. Sci. Technol.* **21** 1592
- [12] Koehler A D et al Atomic layer epitaxy AlN for enhanced AlGaIn/GaN HEMT passivation *IEEE Electron Device Lett.* **34** 1115–7
- [13] Kim D et al 2007 ALD Al₂O₃ passivated MBE-grown AlGaIn/GaN HEMTs on 6H-SiC *Electron. Lett.* **43** 129–30
- [14] Wang H, Chung J W, Gao X, Guo S and Palacios T 2010 Al₂O₃ passivated InAlN/GaN HEMTs on SiC substrate with record current density and transconductance *Phys. Status Solidi C* **7** 2440–4
- [15] Chung J W, Saadat O I, Tirado J M, Gao X, Guo S and Palacios T 2009 Gate-recessed InAlN/GaN HEMTs on SiC substrate Al₂O₃ Passivation *IEEE Electron Device Lett.* **30** 904–6
- [16] Hashizume T, Nishiguchi K, Kaneki S, Kuzmik J and Yatabe Z 2018 State of the art on gate insulation and surface passivation for GaN-based power HEMTs *Mater. Sci. Semicond. Process.* **78** 85–95
- [17] Ghobadi A, Ghobadi T G U, Karadas F and Ozbay E 2018 Angstrom thick ZnO passivation layer to improve the photoelectrochemical water splitting performance of a TiO₂ nanowire photoanode: the role of deposition temperature *Sci. Rep.* **8** 16322
- [18] Li T C et al Surface passivation of nanoporous TiO₂ via atomic layer deposition of ZrO₂ for solid-state dye-sensitized solar cell applications *J. Phys. Chem. C* **113** 18385–90
- [19] Pascoe A R, Bourgeois L, Duffy N W, Xiang W and Cheng Y-B 2013 Surface state recombination and passivation in nanocrystalline TiO₂ dye-sensitized solar cells *J. Phys. Chem. C* **117** 25118–26
- [20] Chandiran A K et al 2012 Subnanometer Ga₂O₃ tunnelling layer by atomic layer deposition to achieve 1.1 V open-circuit potential in dye-sensitized solar cells *Nano Lett.* **12** 3941–7
- [21] Schroder D K 2006 *Semiconductor Material and Device Characterization*. (Hoboken, NJ, USA: Wiley)
- [22] Chen Y, Ma X, Chen W, Hou B, Zhang J and Hao Y 2015 Influence of the gate edge on the reverse leakage current of AlGaIn/GaN HEMTs *AIP Adv.* **5** 097154
- [23] Lee D S et al 2012 Impact of Al₂O₃ Passivation Thickness in Highly Scaled GaN HEMTs *IEEE Electron Device Lett.* **33** 976–8
- [24] Stoklas R et al 2017 Influence of oxygen-plasma treatment on AlGaIn/GaN metal-oxide-semiconductor heterostructure field-effect transistors with HfO₂ by atomic layer deposition: leakage current and density of states reduction *Semicond. Sci. Technol.* **32** 045018
- [25] Huang R et al 2018 Angular dependent XPS study of surface band bending on Ga-polar n-GaN *Appl. Surf. Sci.* **440** 637–42