

Experimental and theoretical investigation of phosphorus in-situ doping of germanium epitaxial layers



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ABSTRACT

We investigate phosphorus in-situ doping characteristics in germanium (Ge) during epitaxial growth by spreading resistance profiling analysis. In addition, we present an accurate model for the kinetics of the diffusion in the in-situ process, modeling combined growth and diffusion events. The activation energy and pre-exponential factor for phosphorus (P) diffusion are determined to be 1.91 eV and $3.75 \times 10^{-5} \text{ cm}^2/\text{s}$. These results show that P in-situ doping diffusivity is low enough to form shallow junctions for high performance Ge devices.

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1. Introduction

Compared to silicon (Si), higher carrier mobility and carrier injection velocity of germanium (Ge) makes it an attractive candidate for electronic devices. In addition, Ge offers a wider absorption window, specially, toward the near infrared region, where Si shows weak or no optical absorption at all. The formation of p-type shallow junctions in Ge (films/layers/substrates) has been successfully demonstrated via ion-implantation technique and in-situ doping during epitaxial growth for boron (B) [1,2] with concentrations above $2 \times 10^{18} \text{ cm}^{-3}$, the solid solubility of B in Ge. However, in the case of n-type Ge junction formation, significant effort is still underway to realize high concentrations and shallow junctions. Arsenic (As) is known to have low solid solubility and high

activation energy (2.71 eV) [3] because its large radius makes it difficult to substitute for Ge atoms [4,5]. In addition, since its high diffusion coefficient ($32 \text{ cm}^2/\text{s}$) [4] hinders the formation of n⁺ shallow junctions in Ge, phosphorus (P), which has relatively high solubility, low activation energy, and low diffusivity, has been widely used. Ion implantation of P can be used to achieve chemical concentrations [6] high 10^{19} cm^{-3} regime, but implant damages [7,8] induce transient enhanced diffusion (TED) hindering formation of shallow junctions in addition to dose losses [6–9], caused by out-diffusion. Doping by conventional furnace diffusion of P leads to deep junctions due to long in-diffusion [5]. In contrast to these doping approaches, in-situ doping technique avoids any implantation-related damages and deep in-diffusions, and consequently provides fairly shallow junctions while offering high electrically activated concentrations [10,11] at low processing temperatures. We recently demonstrated an abrupt and box-shaped high quality Ge n⁺/p junction diode with in-situ doping technique the result of which is reported elsewhere [12]. In order to estimate accurately junction depth and shape when forming n-type shallow junctions in Ge, accurate dopant diffusion modeling is required. Although there are a few theoretical studies about P diffusion in Ge [14,15], by simulations, such analyses are performed

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under the assumption that P atoms are doped by ion implantation method and a simulation model related to P in-situ doping technique is not well established yet. Therefore, in this letter, we investigate the electrical characteristics of P in-situ doped epitaxial Ge layers and follow with a discussion of theoretical diffusion model that can accurately predict our experimental findings.

2. Experiments

After the standard Si wafer cleaning process, a p-type (100) Si wafer was immediately loaded into a cold wall Applied Materials Centura RP-CVD epitaxial reactor. A hydrogen bake at 900 °C was carried out to make sure that no native oxide remained on the Si surface because native oxide on Si wafer prevents Ge film growth on Si substrate. A very thin Si epi layer was first grown using dichlorosilane (DCS) at 700 °C to improve the final film quality. This initial growth was followed by annealing in H₂ ambient for 30 min at 825 °C. Next two Ge buffer growth steps were performed at 400 °C with the growth rate of ~30 nm/min, also followed by anneal in H₂ ambient at 825 °C. The hydrogen annealing process reduced Ge surface roughness by re-flowing the Ge atoms [16]. Then, we performed the final growth step for intrinsic Ge layer at 600 °C with the growth rate of ~60 nm/min. This intrinsic Ge epitaxial layer showed p-type 1 × 10¹⁴ cm⁻³ of electrically activated concentration. Before growing the doped layer, the dopant gas line is purged with diluted 1% phosphine to avoid the dopant cross-contamination at the stabilization step. The final Ge layer was grown at 400 °C–600 °C at 8 Pa on the intrinsic Ge layer, in-situ doped with diluted 1% phosphine for 2 min to form a n⁺ Ge layer.

3. Results and discussion

On an intrinsic Ge virtual substrate, in-situ doping of P is performed by PH₃ and GeH₄ flow for 2 min at the maximum activated doping condition. The mass flow ratios of F(PH₃)/F(GeH₄) are 0.007, 0.009, and 0.01 at the growth temperatures of 400 °C, 500 °C, and 600 °C respectively. Spreading resistance profiling (SRP) measurements on grown samples are obtained. Extracted concentration of P atoms versus depth from the surface is plotted in Fig. 1 at the

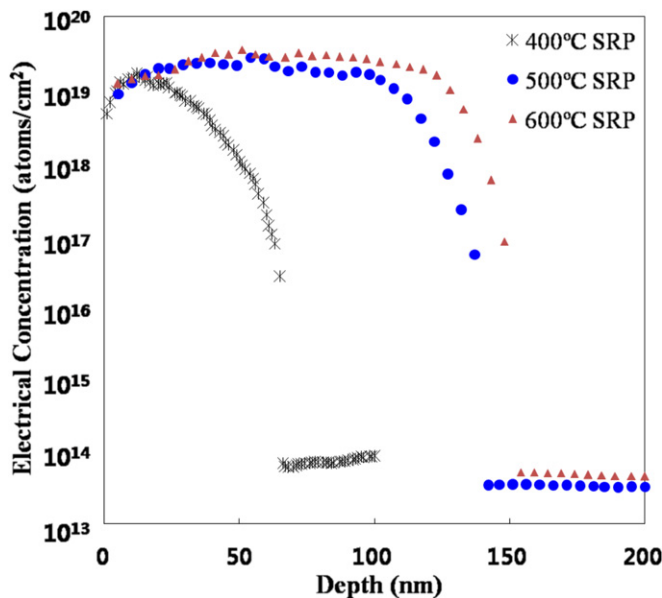


Fig. 1. SRP depth profiles of P in 400 °C, 500 °C, and 600 °C in-situ doped Ge layer growth for 2 min.

growth temperatures of 400 °C–600 °C. Deeper diffusion of P atoms was observed as the deposition temperature increases. The profile exhibits an abrupt and box-shaped junction formation with a slope for the decay of P concentration, 13.7 nm/decade. In addition, the electrically activated dopant profile near the surface have positive slopes such that the electrically activated dopant profiles near the surface is not at maximum because of the shorter anneal time of the deposited phosphorus atoms in this region compared to inner region.

The existing models in literatures [3,6,7,14,15] predicting the diffusion kinetics of P in Ge use error-function profiles. In such models, diffusion is from either an initial dopant profile by ion implantation or assuming a constant dose source at the surface. However, the actual epitaxial growth and in-situ doping process is modeled as consecutive discrete events. Fig. 2 shows the proposed growth model for in-situ doping process. During the in-situ doping process, dopants are provided at the growing surface, and at the same time, diffuse into the substrate. In order to take the effect of a moving surface frontier into account, the growth process is modeled by discrete consecutive cycles. Each cycle is composed of growth phase and diffusion phase. During the growth phase, highly doped thin Ge layer of thickness Δx is assumed to be deposited on the surface. Dopants are diffused during the diffusion phase, and the new doping profile after a short time interval Δt is obtained. A single cycle models in-situ doping growth and dopant diffusion during the short time interval Δt, so growth thickness Δx should be v · Δt, where v is the growth rate of Ge layer (nm/sec). To model the dopant diffusion during the time interval Δt, first, we divide the Ge substrate into very thin layers. If the thickness of the each layer is thin enough, then the doping profile in each layer can be considered as a delta function.

Fick's Second Law is applied to numerically calculate the resulting diffuse profile. The solution of a cycle is used as the initial condition for the subsequent cycle in time.

$$\frac{\partial C}{\partial t} = \frac{\partial}{\partial x} \left(D \frac{\partial C}{\partial x} \right) \tag{1}$$

Due to high doping concentrations required in device applications, concentration-dependent-diffusion and electric-field (E-field) effects were included in our model. The following relations were used for second order concentration dependency and E-field effects:

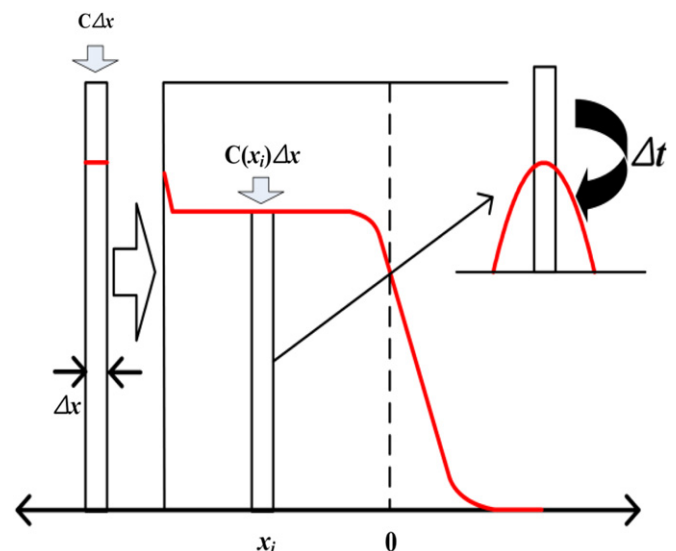


Fig. 2. Schematic of in-situ growth model.

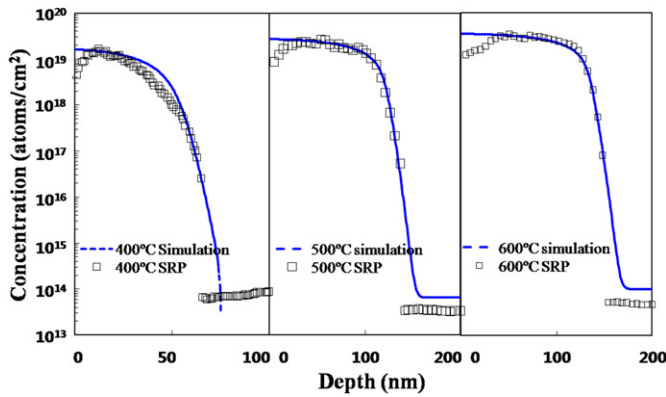


Fig. 3. Phosphorus (P) concentration profiles of experimental and simulated data at 400 °C, 500 °C and 600 °C.

$$D = D^0 + D^- \left(\frac{n}{n_i} \right) + D^{--} \left(\frac{n}{n_i} \right)^2 \quad (2)$$

$$\frac{\partial C}{\partial t} = \frac{\partial}{\partial x} \left(D \left[\frac{\partial C}{\partial t} \right] \right) + \frac{\partial}{\partial x} \left(DC \frac{\partial}{\partial x} \left[\ln \left(\frac{n}{n_i} \right) \right] \right) \quad (3)$$

where the first term $\partial/\partial x(D[\partial C/\partial t])$ is related to basic Fick's Second Law and the second term $\partial/\partial x(DC\partial/\partial x[\ln(n/n_i)])$ is related to the E-field effect. In order to solve these equations, the derived solution in Asaithambi [13] and finite-difference method is used.

The concentrations of species obtained by SRP data are electrically active concentrations. It was shown that the electrical activation levels for 400 °C, 500 °C, and 600 °C in-situ P-doped Ge epitaxy layers were 30%, 65%, and 100% respectively [12]. In the model described, the chemical concentration of P in Ge was used by scaling the measured electrically active concentration by the level of activation at each growth temperature. As shown in Fig. 3, the resulting profiles calculated numerically with our model and the experimental results are in very good agreement at 400 °C, 500 °C, and 600 °C respectively.

At a given growth temperature, the doubly charged diffusivity, D^{--} , dominates due to the quadratic dependence on the carrier concentration. This supports the presence of concentration

dependent effects in addition to a sharp decaying tail of the measured profile. The junction depth, therefore, is governed by the doubly charged diffusivity of P atoms. Fig. 4 plots the extracted doubly charged diffusivity coefficient and corresponding Arrhenius curve to fit experimental data. It is shown that the activation energy and pre-exponential factor for phosphorus diffusion are determined to 1.91 eV and $3.75 \times 10^{-5} \text{ cm}^2/\text{s}$. The activation energy corresponding to the doubly charged diffusivity, D^{--} , in this work is slightly lower than that observed in Chui et al. (2.07 eV) [5], Carroll et al. (2.3 eV) [14] and Tsouroutas et al. (2.69 eV) [15]. Chui et al. and Carroll et al. investigated the concentration dependent diffusivity of P in Ge when the dopants are introduced by ion-implantation technique. The pre-factor of the second order diffusivity coefficient obtained in this work is dramatically lower than that reported by Chui et al. ($4.38 \times 10^{-2} \text{ cm}^2/\text{s}$) and Carroll et al. ($1.85 \times 10^{-2} \text{ cm}^2/\text{s}$). Therefore, in-situ doping technique can offer the low diffusivity unattainable by earlier work [5,14], paving the way for fabricating shallow n-type junctions and devices with high levels of electrically active P in Ge.

4. Conclusion

We investigated characteristics of P in-situ doped Ge obtained by epitaxial growth. The accurate modeling of P in-situ doping process was achieved with the activation energy of 1.91 eV and pre-exponential factor of $3.75 \times 10^{-5} \text{ cm}^2/\text{s}$. Phosphorus in-situ doping diffusivity is low enough to form shallow junction devices with high performance in Ge substrate.

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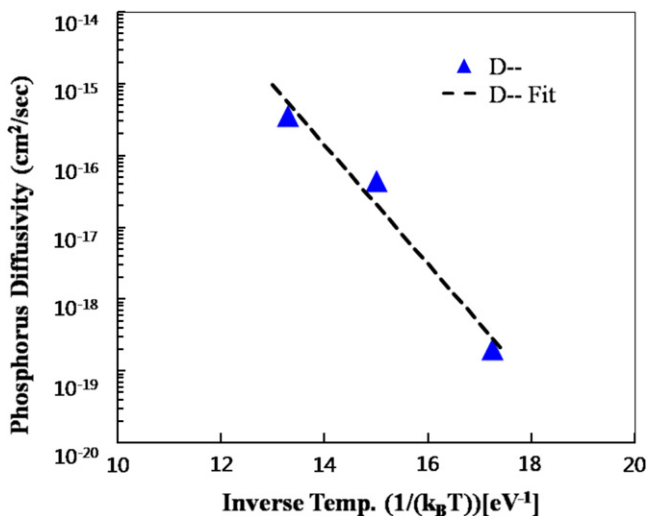


Fig. 4. Plot of diffusivity coefficients (D^{--}) at different temperatures. For D^{--} , activation energy is 1.91 eV, and diffusivity pre-factor is $3.75 \times 10^{-5} \text{ cm}^2/\text{s}$.

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