

Chapter 2

Colloidal Semiconductor Nanocrystals



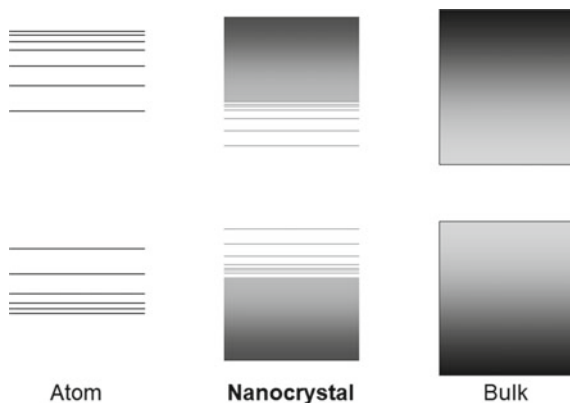
Abstract In this chapter, we review colloidal semiconductor nanocrystals (NCs) and their remarkable size-dependent properties. We emphasize on colloidal nanoplatelets and explain how they differ from NCs of other classes.

Keywords Semiconductor · Colloidal nanocrystals · Quantum dots · Quasi-2D nanoplatelets

The presence of man-made nanocrystals (NCs) can be traced back to as early as three millennia ago [1]. However, compared to this time scale, the achievement of their reproducible synthesis using colloidal chemistry three decades ago is quite recent [2]. This rather recent development, though, enabled an enormous leap in semiconductor optoelectronics ever since. After nearly 30 years of their first well-controlled synthetic introduction, intensive research keeps being conducted on studying, improving, and exploiting their efficiency and other properties, as well as their incorporation in optoelectrical applications.

The magnificence of NCs lies in their size-dependent physical properties. This size dependence can be best understood by the quantum confinement effect, which can be quantified to first order by treating the NC as a quantum well with a finite barrier potential. Even by ignoring the crystalline structure and the resulting periodic potential within the NC volume, this treatment is able to approximate the evolution of electronic states with the dimensions of the quantum well [3]. In general, quantum confinement starts to become observable when the dimensions of the NC are comparable to the exciton Bohr radius. As the size of the NC keeps shrinking, this will result in further separation of the energy levels of the system. In the case of semiconductors, the separation between the highest occupied energy band and the lowest unoccupied energy band, i.e. the bandgap, will increase with shrinking size. This has been schematically demonstrated in Fig. 2.1. Not only does the bandgap increase with reduced size but also the energy bands are discretized with increased separation between states as the NC shrinks. In addition to the modifications in the energy levels, uncertainty will be introduced in the momenta of electrons and holes due to

Fig. 2.1 Energy states of an atom, bulk crystal, and nanocrystal. Nanocrystals have discrete energy states near the highest occupied and lowest unoccupied molecular orbitals



their being localized within the NC volume, resulting in relaxation of the momentum conservation [4].

The work that pioneered the modern colloidal synthesis of semiconductor NCs came in 1993 by Murray et al. [2]. Herein, the formation of CdSe, CdS, or CdTe nanoparticles in colloidal medium is initiated by injection of Cd and Se (or S, Te) precursors into solution at an elevated temperature; these nuclei are grown until the desired size is reached and the growth is terminated by injecting organic ligands and lowering the reaction temperature. The resulting ensembles of CdSe spherical quantum dots (QDs) in this work have size dispersity as low as 5%, with the average QD size being tunable from 1.2 to 11.5 nm. The absorption onset of these CdSe QD ensembles is ~ 400 nm for the smallest average size and ~ 700 nm for the largest one [2].

An exemplary transmission electron micrograph of QDs is displayed in Fig. 2.2a. The crystallographic planes of the individual QDs are visible in the electron microscopy image. The crystalline structure is also depicted schematically in Fig. 2.2b. Typically, surface planes of these NCs are saturated by long-chain organic ligands. These organic ligands are used as a means to render the NC ensembles soluble in nonpolar organic solvents such as hexane and toluene. Additionally, they help with the passivation of the highly energetic, unsaturated bonds of the atoms on the NC surface. This kind of unsaturated bonds can be detrimental to the optical efficiency of NCs by resulting in charge trapping at the surface, which reduces the photoluminescence (PL) efficiency and the charge mobility, and therefore is undesirable in optical and photovoltaic applications [5]. Since the surface-to-volume ratio of the NCs is greatly enhanced in these nanometer-sized particles in comparison with bulk materials, loss of charge carriers and emission through surface trapping can significantly deteriorate the efficiency of these QDs, and saturation of the atomic bonds on the surface is essential.

The most common ligands are organic molecules with long hydrocarbon chains attached to a functional group. Examples of such ligands used in the NC synthesis include oleic acid [6, 7], trioctylphosphine [8, 9], and trioctylphosphine oxide [10].

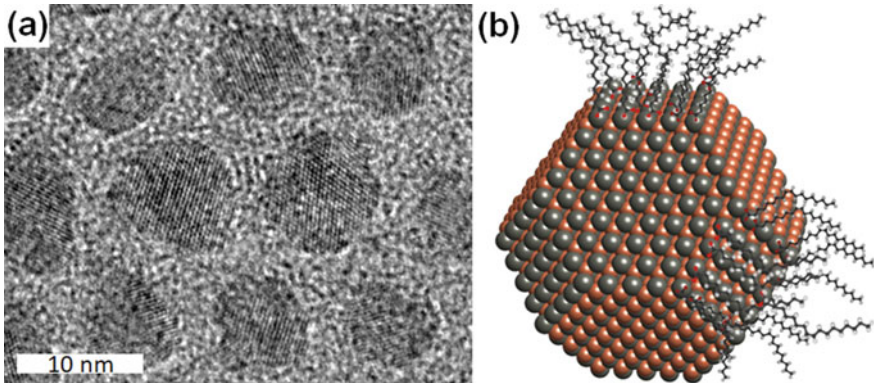


Fig. 2.2 **a** Transmission electron micrograph of CdZnS/ZnS quantum dots (2018 Demir Group). **b** Schematic depiction of a quasi-spherical colloidal CdSe nanocrystal capped with ligands (in this schematic, oleate) on facets. Ligands on some facets are not drawn for clarity purposes

In most cases, the functional group of these molecules is ionized in solution and attaches itself to the oppositely charged atoms on the surface [11].

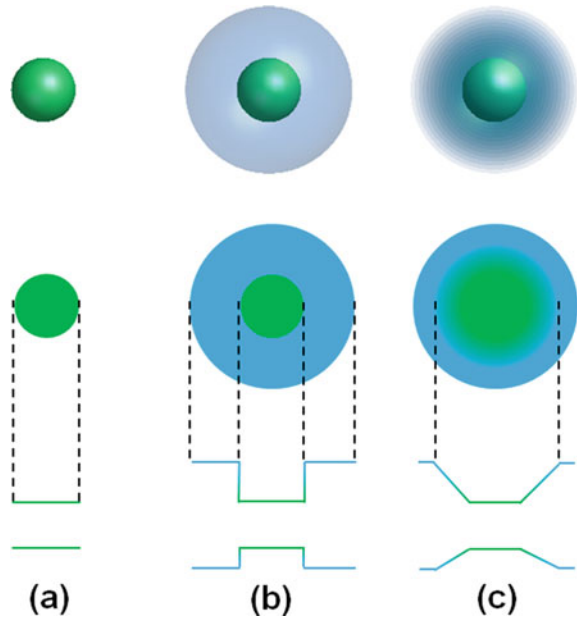
It is also possible to bind inorganic ligands to the surface of QDs to achieve solubility in polar solvents. Approaches for QD synthesis in water-based environment [12] or post-synthesis procedures for replacing organic ligands with inorganic ones enable the making of water-soluble NCs [13, 14].

Today, synthesis of highly monodisperse QD ensembles with high PL efficiency and narrow emission linewidth is possible with the emission wavelength tunable through the entire visible spectrum and beyond. To this end, a variety of compositions, as well as heterostructures have been introduced. While Cd-based II–VI QDs are still being studied and used to a great extent, other compositions have been eventually introduced not long after CdX QDs. These compositions include III–V semiconductors such as InAs [9] and InP [15, 16] and IV–VI semiconductors such as PbS [6] and PbSe [17]. Owing to their much smaller bandgap than CdSe, the QDs of these other semiconductors are more suitable for the IR regime. Altogether, the spectral range of QDs made of these material systems can extend from visible to IR (up to 3.5 μm) [18, 19].

Even more common than the single-composition QDs (Fig. 2.3a) are the QD heterostructures, where the “core” QD is enclosed by a “shell” to increase the surface passivation and prevention of leakage of the excitonic wavefunction outside the QD volume (Fig. 2.3b). With such heterostructures as CdSe/CdS [20, 21], CdSe/ZnS [22, 23] and InP/ZnS [24, 25] more stable solutions and better quantum yield is achieved compared with core-only QDs. It is also possible to create a potential gradient between the core and shell, where the molecular composition changes gradually from the center to the outermost shell (Fig. 2.3c) [26–28].

Colloidal QDs have advantageous properties including but not limited to emission/absorption tunability, solution processability, and large absorption cross-section. Furthermore, as their sizes are orders of magnitude smaller than light wavelength,

Fig. 2.3 Two- and three-dimensional schematics for **a** core, **b** core/shell, and **c** gradient core/shell QDs. Depending on the selected pair of compositions, the valence and conduction band offsets of core and (gradient) shell may vary



they do not induce scattering individually while interacting with light. These properties render them especially favorable in optoelectrical devices. Applications including light-emitting diodes (LEDs) [29], lasers [10, 30], and field effect transistors [31] have already benefited from the favorable features of these tunable nanoemitters.

Readers interested in a detailed overview of the synthesis and applications of various NCs are referred to [19, 32]. A theoretical treatment of semiconductor NCs can be found in [3, 33].

2.1 Colloidal Quantum Wells (Nanoplatelets)

Apart from the deeply investigated spherical QDs, colloidal NCs of various shapes and dimensionalities have also been eventually reported and studied. These include nanocubes [34], tetrapods [35, 36] and quasi-one-dimensional (1D) nanorods [37, 38]. In the last decade, a new class of colloidal NCs has been introduced, which is the quasi-two-dimensional nanoplatelets (NPLs) with one-dimensional quantum confinement [39, 40].

Colloidal NPLs have lateral sizes ranging from several nm to about 100 nm while having a thickness of only a few atoms. This means that, unlike spherical, quasi-0D QDs, where the quantum confinement is 3D, NPLs have effective quantum confinement only along the vertical direction. The thickness of the NPL therefore determines the strength of quantum confinement, and, in turn, the bandgap.

The most remarkable property of the NPLs is that their lateral surfaces are atomically flat with no variation in thickness. Owing to this atomic precision in their thickness, their ensembles have little or no inhomogeneous broadening. NPLs can thus have emission linewidths as small as several nms [39] at room temperature, which has not been possible with QDs even in their highly monodisperse ensembles.

The presence of NPLs has been observed in the early 2000s as ultra-narrow features in absorption and emission features of colloidal NC ensembles. Initially, these features have been attributed to “magic-sized” nanoclusters in the ensemble with discrete sizing [41, 42]. At the end of the decade, Ithurria et al. managed to reproducibly synthesize these discretely sized CdX ($X = \text{Se, S, Te}$) NCs, not as a by-product, but as the main output of their synthesis procedure. Their seminal work demonstrated that NPLs are atomically flat with one-dimensional quantum confinement, giant oscillator strength, and ultranarrow emission linewidth [39, 40].

In Fig. 2.4a, individual CdSe NPLs of different thicknesses are drawn schematically. The crystal structure of a NPL is composed of alternating atomic planes of Cd and Se, the bottom- and topmost layers always being a Cd plane. Thus, for N layers of Se planes in an NPL, there will be $N + 1$ Cd planes. In this case, the thickness of the NPL can be expressed as N monolayers (ML), where a “monolayer” represents a Cd-Se atomic pair along the vertical direction. It is also not uncommon to use half-integers as the number of monolayers, to signify the presence of the extra Cd atomic plane. Throughout this brief, we use this latter nomenclature when referring to the core NPL thicknesses.

Figure 2.4b displays the scanning transmission electron micrograph of 5.5 ML CdSe NPLs. For most of the NPLs in the image, the visible surface is a lateral one. In Fig. 2.4c, absorption and photoluminescence spectra of 3.5, 4.5, and 5.5 ML NPL ensembles are plotted. The sharp excitonic features can be discerned in these spectra, which correspond to the electron-heavy hole and electron-light hole transitions. The Stokes shift between the emission and absorption peaks is exceptionally low, typically limited to a few nm [39]. It is also seen here that the atomically flat thickness and the confinement energy being essentially determined by only one dimension comes at the cost of the discretization of spectral tunability. In contrast to QDs, where the excitonic features can be continuously tuned by size engineering, the size-dependent absorption peak leaps from ~ 460 nm to 510 nm when the NPL thickness is increased from 3.5 to 4.5 ML, and to 553 nm when it is further increased by another ML. This limitation in tunability can be largely overcome by various approaches such as compositional alloying [43, 44], making heterostructures and to a limited extent, enforcing additional quantum confinement along one of the two lateral dimensions, i.e. making the size along that lateral dimension comparable to the exciton Bohr radius [45, 46].

To date, the synthesis of NPLs with thicknesses ranging from 2.5 to 9.5 ML has been reported [39, 46–51]. In addition, similar to QDs, making heterostructures of NPLs is also possible for purposes of improved colloidal stability and photoluminescence efficiency. However, unlike QDs, where the only possibility is growing a radial shell, there are two degrees of freedom when shelling NPLs, namely the growth of lateral wings around the periphery of the NPL, and sandwiching the core NPL by

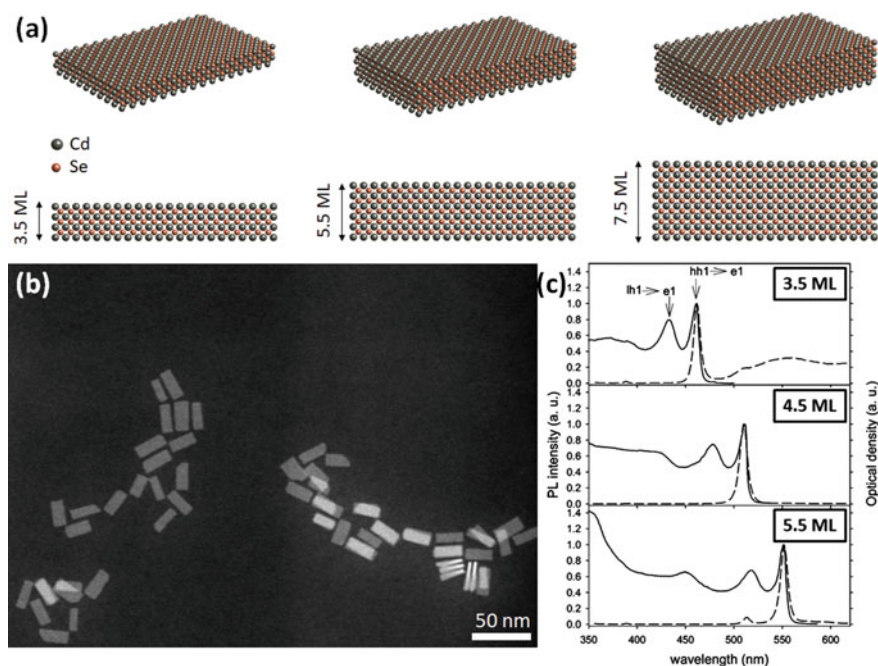
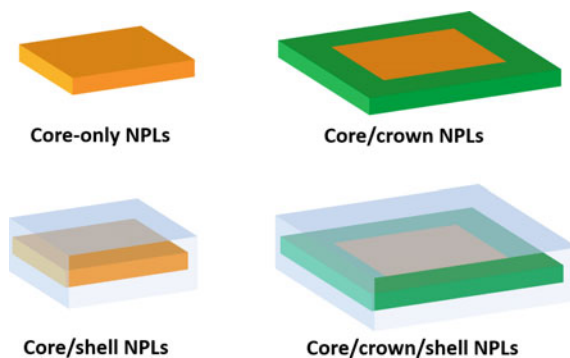


Fig. 2.4 **a** Schematic drawing of 3.5, 5.5 and 7.5 ML CdSe zinc-blend nanoplatelets (NPLs). **b** Scanning transmission electron micrograph of 5.5 ML NPLs (2017 Demir Group). **c** Absorbance (solid line) and PL (dashed line) spectra of 3.5, 4.5, and 5.5 ML CdSe NPLs. Adapted with permission from [40]. Copyright 2008 American Chemical Society

shelling it vertically and surrounding a little bit around the periphery as well. These heterostructures, commonly known as core/crown [52, 53] and core/shell NPLs, respectively, are schematically depicted in Fig. 2.5. It is also possible to combine these two shelling methods to obtain core/crown/shell NPLs [54].

Fig. 2.5 Schematics of core, core/shell, core/crown, and core/crown/shell NPLs



NPLs have already been demonstrated to be suitable for optoelectronics applications such as LEDs [43, 55], solar concentration [56], and lasing [57–59], in most cases outperforming their quantum-dot counterparts. This superior performance is thanks to their enhanced absorption, giant oscillator strength, and lack of inhomogeneous broadening.

Another striking property of the NPLs is the anisotropy of the band-edge excitonic dipole, which leads to directional emission out of NPLs [60]. Such directionality is essential for the active media of LEDs and lasers and is one of the main motivations for controlling the orientation of NPLs collectively in their ensemble. Methods of orientation control of NPLs and their use in optoelectronic applications will be discussed in detail in Chap. 4 and 5.

For a more detailed overview of colloidal NPLs, a number of reviews are available that cover early and more recent progress on synthesis, optical properties, and potential applications [32, 53, 61, 62].

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